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Gravitational radiation from inspiralling compact objects: Spin effects to fourth Post-Newtonian order

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The linear- and quadratic-in-spin contributions to the binding potential and gravitational-wave flux from binary systems are derived to next-to-next-to-leading order in the Post-Newtonian (PN) expansion of general relativity, including finite-size and tail effects. The calculation is carried out through the worldline effective field theory framework. We find agreement in the overlap with the available PN literature and test-body limit. As a direct application, we complete the knowledge of spin effects in the evolution of the orbital phase for aligned-spin circular orbits to fourth PN order. We estimate the impact of the new results in the number of accumulated gravitational-wave cycles. We find they will play an important role in providing reliable physical interpretation of gravitational-wave signals from spinning binaries with future detectors such as LISA and the Einstein Telescope.

Introduction. The dynamical evolution of compact binaries has been the main cause of the gravitational waves (GWs) detected by the LIGO-Virgo-KAGRA interferometers [1–3], and will continue to be one of the primary sources for future GW observatories such as the Laser Interferometer Space Antenna (LISA) [4] and the Einstein Telescope (ET) [5]. The GWs produced by the inspiral, merger, and ring-down from the expected several two-body events will carry vast amounts of information that can shed light on long-standing problems in astrophysics, cosmology, and particle physics [6–9]. In particular the spin of the constituents, which has been found to be large in several recent detections [3], is not only strongly correlated with different formation channels, e.g. [3, 10–15], also offers a window to physics beyond the standard model, e.g. [16–24]. Therefore, high-precision waveforms incorporating spin corrections are an essential ingredient to exploit the discovery potential in GW astronomy.

After a concerted effort involving both numerical [25–28] as well as analytic techniques [29–31], GW template banks have been successfully used to analyse the GW data collected thus far [1–3]. However, while current templates may be sufficient for detection, when it comes to parameter estimation, the formidable empirical reach of future experiments require higher levels of accuracy, both for the Post-Newtonian (PN) inspiral regime as well as merger stages of the binary’s dynamics [32–35]. Presently, although partial results for the derivation of the evolution of the orbital phase in the inspiral regime are known at 4PN order for non-spinning bodies through various computations [36–48], and even higher orders in the conservative sector [49–56], spin contributions have not been pushed so far to the same relative level of accuracy. In particular, for radiative effects, while spin-orbit corrections were obtained to next-to-next-to-leading order (N²LO) [57, 58], complete spin-spin effects are only known to NLO [59–66]. In this letter we fill this gap and report the completion of spin effects to N²LO in the PN expansion and to quadratic order in the spins, corresponding to the 4PN order for rapidly rotating bodies.

The derivation involves several ingredients, which we obtain using the worldline effective field theory (EFT) framework in the PN regime [67, 68] extended to spinning bodies [59–64, 69]. The EFT approach uses powerful tools from particle physics resembling, for instance, methodologies used in the calculation of binding energies for heavy quark states [70]. The problem of motion is thus reduced to a series of Feynman diagrams, involving potential and radiation modes, which are constructed by iteratively solving for the (classical) gravitational field sourced by compact objects treated as point-like objects. Utilizing the EFT formalism, we have compute the gravitational potential and necessary radiative multipole moments at linear and quadratic order in the spins entering in the flux to N²LO, including finite-size effects. The completion of spin contributions at 4PN entails also the so-called *tail effect*, due to the scattering of the outgoing radiation off of the background geometry, e.g. [42], which we incorporate through the EFT approach. The values for all of the intermediate (very lengthy) results are displayed in the ancillary file, see also the supplemental material. Perfect agreement is found in the overlap with previous results in the literature [57, 58, 71–77].

From the binding energy and radiated flux we derive the imprint of spin effects in the orbital phase evolution for aligned-spins circular orbits. As a measure of the impact of the new terms, we estimate the accumulated GW cycles for various paradigmatic astrophysical configurations as well as detector-sensitivity curves. We find the N²LO spin terms make a significant contribution both in ET and LISA frequency bands. (The effect increases the larger the mass ratio, which coincidentally is expected to correlate with larger spins [13].) This is partially driven by quadratic-in-spin terms carrying information about the inner structure of the compact objects. The results presented here will therefore play an important role in elucidating the origin of binary black holes as well as aiding future discoveries, such as new dark objects [24] or clouds of putative ultralight particles induced by superradiance [16–23], through GW precision data.

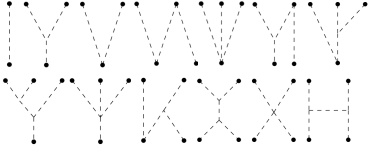


Figure 1. Topologies needed for the N^2 LO potential (see text).

Worldline EFT approach. The effective action is obtained in the weak-field regime by *solving for* the metric perturbation in the (classical) saddle-point approximation [59, 67]. The compact bodies are described by the *Routhian*,

$$\mathcal{R} = -\frac{1}{2} \sum_{n=1,2} \left(m_n g_{\mu\nu} v_n^\mu v_n^\nu + \omega_\mu^{ab} S_{n ab} v_n^\mu \right. \quad (1)$$

$$\left. - \frac{C_{\text{ES}^2}^{(n)}}{m_n} \frac{E_{ab} S_n^{ac} S_n^b}{\sqrt{v_n^\mu v_{n\mu}}} + \frac{1}{m_n} R_{deab} S_n^{ab} S_n^{cd} v_n^e v_{n c} + \dots \right),$$

which serves both as a Lagrangian for the position variables (x_n^α, v_n^α) and a Hamiltonian for the spin, projected onto a locally-flat frame, $S_n^{ab} \equiv e_\mu^a e_\nu^b S_n^{\mu\nu}$, and coupled to ω_μ^{ab} , the Ricci rotation coefficients. The free parameters include the masses, m_n , as well as $C_{\text{ES}^2}^{(n)}$, which accounts for finite-size effects [59–62]. The latter couple to E_{ab} , the electric component of the Weyl tensor. The last term in (1), involving the Riemann tensor, ensures the covariant spin-supplementarity-condition, $S^{ab} v_b = 0$, is preserved upon evolution. However, for convenience, our results will be presented in terms of (Newton-Wigner) precession-only spin variables. The ellipses encapsulate higher orders in spin and curvature, which are not relevant in this letter. See [31] for more details.

Gravitational potential & Binding energy. The derivation of the potential follows by computing the ‘vacuum-to-vacuum’ amplitude in the presence of external sources, by *integrating out* the off-shell quasi-instantaneous modes of the gravitational field. The associated Feynman diagrams needed to N^2 LO are depicted in Fig. 1. The worldline couplings, depicted as a black disc, include the mass as well as the linear and quadratic spin terms shown in (1). The dashed lines represent the potential modes of the gravitational field responsible for the binding of the two-body system. Hence, the Green’s function (a.k.a. propagator) must be PN expanded [67]

$$\frac{i}{(k^0)^2 - \mathbf{k}^2} = -\frac{i}{\mathbf{k}^2} \left(1 + \frac{(k^0)^2}{\mathbf{k}^2} + \dots \right), \quad (2)$$

with each factor of $(k^0)^2$ scaling as v^2 .

From the gravitational potential, and after transforming to conserved-norm spins, we derive the equations of motion including the spin precession, which are displayed in full glory in the ancillary file. We then compute the

binding energy, which can be written as

$$E_{\text{SO}} = \frac{GM\nu}{r^2} \left[e_3^0 + \frac{1}{2} \left(e_5^0 + \frac{GM}{r} e_5^1 \right) \right] + \frac{GM\nu}{4r^2} \left[\frac{1}{2} \left(e_7^0 + \frac{GM}{r} e_7^1 \right) + \frac{G^2 M^2}{r^2} e_7^2 \right], \quad (3)$$

for the spin-orbit terms, and

$$E_{\text{SS}} = \frac{G\nu}{4r^3} \left[e_4^0 + \frac{1}{2} e_6^0 + \frac{GM}{r} e_6^1 \right] + \frac{G\nu}{16r^3} \left[\frac{1}{2} e_8^0 + \frac{GM}{r} e_8^1 + \frac{7}{2} \frac{G^2 M^2}{r^2} e_8^2 \right], \quad (4)$$

for spin-spin contribution. The value for the e_j^i PN coefficients are given in the supplemental material and ancillary file. We use $M \equiv m_1 + m_2$ for the total mass, $\nu \equiv m_1 m_2 / M^2$ for the symmetric-mass-ratio and $r \equiv |\mathbf{x}_1 - \mathbf{x}_2|$ the relative distance, respectively. We use the parameter $\delta \equiv (m_1 - m_2) / M$, and $\kappa_\pm \equiv C_{\text{ES}^2}^{(1)} \pm C_{\text{ES}^2}^{(2)}$. For the spin variables we use the standard, e.g. [29],

$$\mathbf{S} \equiv \mathbf{S}_1 + \mathbf{S}_2, \quad (5)$$

$$\boldsymbol{\Sigma} \equiv M \left(\frac{\mathbf{S}_2}{m_2} - \frac{\mathbf{S}_1}{m_1} \right).$$

Radiated flux. The emitted power is obtained by matching the one-point function to a long-distance worldline effective theory for the binary system treated as a point-like object endowed with multipole moments. The action includes, in addition to the (Bondi) mass-energy monopole term, $-M_B \int d\tau$, a series of symmetric-trace-free (STF) *source* mass, I^L , and current, J^L , time-dependent multipoles, with $L \equiv \{i_1 \dots i_L\}$, [68]

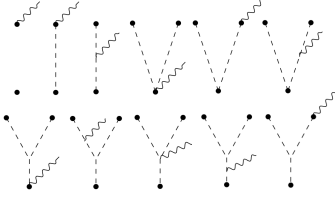
$$\sum_{\ell=2} \left(\frac{1}{\ell!} I_{\text{STF}}^L \nabla_{L-2} \bar{E}_{i_{\ell-1} i_\ell} - \frac{2\ell}{(2\ell+1)!} J_{\text{STF}}^L \nabla_{L-2} \bar{B}_{i_{\ell-1} i_\ell} \right), \quad (6)$$

which couple to (covariant) derivatives of \bar{E}_{ij} and \bar{B}_{kl} , the electric and magnetic components of the Weyl tensor involving only the *background* radiation field. The relevant topologies are shown in Fig. 2. The wavy line represents the on-shell radiation, which couples both to the constituents of the binary as well as the binding modes. From the source multipoles we compute the energy flux by squaring the emission amplitude [68],

$$\mathcal{F}_{\text{src}} = G \left(\frac{1}{5} I_{ij}^{(3)} I_{ij}^{(3)} + \frac{16}{45} J_{ij}^{(3)} J_{ij}^{(3)} + \frac{1}{189} I_{ijk}^{(4)} I_{ijk}^{(4)} \right. \quad (7)$$

$$\left. + \frac{1}{84} J_{ijk}^{(4)} J_{ijk}^{(4)} + \frac{1}{9072} I_{ijkl}^{(5)} I_{ijkl}^{(5)} + \frac{4}{14175} J_{ijkl}^{(5)} J_{ijkl}^{(5)} \dots \right),$$

to the desired order. The time derivatives are computed within the adiabatic approximation, by using the conservative equations of motion.

Figure 2
entier

oments

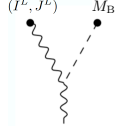


Figure 3. Tail coupling between the binary's mass monopole and other source moments.

We find for the spin-orbit flux

$$\begin{aligned} \mathcal{F}_{\text{src}}^{\text{SO}} = & \frac{8G^3 M^3 \nu^2}{15r^5} \left(\left[f_3^0 + 4 \frac{GM}{r} f_3^1 \right] \right. \\ & + \frac{1}{28} \left[f_5^0 + 2 \frac{GM}{r} f_5^1 + 4 \frac{G^2 M^2}{r^2} f_5^2 \right] \\ & \left. + \frac{1}{84} \left[f_7^0 + \frac{GM}{r} f_7^1 + \frac{2}{3} \frac{G^2 M^2}{r^2} f_7^2 + \frac{4}{9} \frac{G^3 M^3}{r^3} f_7^3 \right] \right), \end{aligned} \quad (8)$$

whereas, for the spin-spin terms, we have

$$\begin{aligned} \mathcal{F}_{\text{src}}^{\text{SS}} = & \frac{2G^3 M^2 \nu^2}{15r^6} \left(f_4^0 + \frac{1}{7} \left[f_6^0 + \frac{GM}{r} f_6^1 + 4 \frac{G^2 M^2}{r^2} f_6^2 \right] \right. \\ & \left. + \frac{1}{84} \left[f_8^0 + \frac{GM}{r} f_8^1 + \frac{2}{3} \frac{G^2 M^2}{r^2} f_8^2 + \frac{8}{3} \frac{G^3 M^3}{r^3} f_8^3 \right] \right). \end{aligned} \quad (9)$$

The value of the f_j^i coefficients are displayed in the supplemental material and ancillary file.

In order to complete the derivation of the total flux we must also include the tail effect, depicted in Fig. 3, where the radiated field interacts with the background geometry around the binary system, sourced by the monopole term. This is often packaged in terms of *radiative* multipole moments, which can then be used to compute the total power using (7). The (leading) tail term yields

$$\begin{aligned} \mathcal{F}_{\text{tail}} = & -G^2 M_B \pi \int \left[\left(\frac{2}{5} I_{ij}(p) I_{ij}(q) + \frac{32}{45} J_{ij}(p) J_{ij}(q) \right) \right. \\ & \left. - pq \left(\frac{2}{189} I_{ijk}(p) I_{ijk}(q) + \frac{1}{42} J_{ijk}(p) J_{ijk}(q) \right) \right] \times \\ & p^3 q^4 \text{sign}(q) e^{-i(p+q)t} dp dq, \end{aligned} \quad (10)$$

where $I_L(p)$ is the Fourier transform. In the above expression the M_B includes not only the total mass of the binary but also the kinetic energy and binding potential to a given PN order.

Aligned-spin circular orbits. As a direct application of our results we consider the phenomenologically relevant case of (planar) circular orbits with the (conserved-norm) spins being either aligned or anti-aligned with the angular momentum. In what follow we quote the results using the following projected (dimensionless) spin variables

$$\hat{S}_\ell \equiv \frac{\boldsymbol{\ell} \cdot \mathbf{S}}{GM^2}, \quad \Sigma_\ell \equiv \frac{\boldsymbol{\ell} \cdot \boldsymbol{\Sigma}}{GM^2}, \quad (11)$$

with $\boldsymbol{\ell}$ the unit vector in the direction of the angular momentum. Garnering all the pieces together we find for the spin-dependent linear- and bilinear-in-spin binding energy as a function of the orbital frequency, $x \equiv (GM\Omega)^{2/3}$,

$$\begin{aligned} E_{\text{spin}} = & -\frac{M\nu x}{2} \left[x^{3/2} \left(\frac{14}{3} \hat{S}_\ell + 2\delta \hat{\Sigma}_\ell \right) + x^2 \left\{ (-2 - \kappa_+) \hat{S}_\ell^2 + \left[\kappa_- + \delta(-2 - \kappa_+) \right] \hat{S}_\ell \hat{\Sigma}_\ell \right. \right. \\ & + \left. \left[2\nu + \frac{1}{2} \delta \kappa_- + \left(-\frac{1}{2} + \nu \right) \kappa_+ \right] \hat{\Sigma}_\ell^2 \right\} + x^{5/2} \left[\left(11 - \frac{61}{9} \nu \right) \hat{S}_\ell + \left(3 - \frac{10}{3} \nu \right) \delta \hat{\Sigma}_\ell \right] \\ & + x^3 \left(\left[\frac{50}{9} + \frac{5}{3} \nu - \frac{5}{3} \delta \kappa_- + \left(-\frac{25}{6} + \frac{5}{6} \nu \right) \kappa_+ \right] \hat{S}_\ell^2 + \left\{ \left(\frac{5}{2} + \frac{35}{6} \nu \right) \kappa_- + \delta \left[\frac{25}{3} + \frac{5}{3} \nu + \left(-\frac{5}{2} + \frac{5}{6} \nu \right) \kappa_+ \right] \right\} \hat{S}_\ell \hat{\Sigma}_\ell \right. \\ & + \left. \left[5 - 10\nu - \frac{5}{3} \nu^2 + \left(\frac{5}{4} + \frac{5}{4} \nu \right) \delta \kappa_- + \left(-\frac{5}{4} + \frac{5}{4} \nu - \frac{5}{6} \nu^2 \right) \kappa_+ \right] \hat{\Sigma}_\ell^2 \right) + x^{7/2} \left[\left(\frac{135}{4} - \frac{367}{4} \nu + \frac{29}{12} \nu^2 \right) \hat{S}_\ell + \left(\frac{27}{4} - 39\nu + \frac{5}{4} \nu^2 \right) \delta \hat{\Sigma}_\ell \right] \\ & + x^4 \left(\left[\frac{67}{12} + \frac{6445}{108} \nu - \frac{7}{36} \nu^2 + \left(-\frac{31}{4} + \frac{35}{18} \nu \right) \delta \kappa_- + \left(-\frac{125}{8} + \frac{1025}{72} \nu - \frac{7}{72} \nu^2 \right) \kappa_+ \right] \hat{S}_\ell^2 + \left\{ \left(\frac{63}{8} + \frac{449}{24} \nu - \frac{553}{72} \nu^2 \right) \kappa_- \right. \right. \\ & + \left. \left. \delta \left[\frac{49}{4} + \frac{1649}{36} \nu - \frac{7}{36} \nu^2 + \left(-\frac{63}{8} + \frac{295}{24} \nu - \frac{7}{72} \nu^2 \right) \kappa_+ \right] \right\} \hat{S}_\ell \hat{\Sigma}_\ell + \left[\frac{21}{2} - \frac{119}{12} \nu - \frac{135}{4} \nu^2 + \frac{7}{36} \nu^3 + \left(\frac{63}{16} + \frac{77}{48} \nu - \frac{91}{48} \nu^2 \right) \delta \kappa_- \right. \right. \\ & \left. \left. + \left(-\frac{63}{16} + \frac{301}{48} \nu - \frac{499}{48} \nu^2 + \frac{7}{72} \nu^3 \right) \kappa_+ \right] \hat{\Sigma}_\ell^2 \right). \end{aligned} \quad (12)$$

On the other hand, the energy flux becomes

$$\begin{aligned}
\mathcal{F}_{\text{spin}} = & \frac{32\nu^2 x^5}{5G} \left[x^{3/2} \left(-4\hat{S}_\ell - \frac{5}{4}\delta\hat{\Sigma}_\ell \right) + x^2 \left\{ (4+2\kappa_+) \hat{S}_\ell^2 + \left[-2\kappa_- + \delta(4+2\kappa_+) \right] \hat{S}_\ell \hat{\Sigma}_\ell \right. \right. \\
& + \left. \left[\frac{1}{16} - 4\nu - \delta\kappa_- + \kappa_+ (1-2\nu) \right] \hat{\Sigma}_\ell^2 \right\} + x^{5/2} \left[\left(-\frac{9}{2} + \frac{272}{9}\nu \right) \hat{S}_\ell + \delta \left(-\frac{13}{16} + \frac{43}{4}\nu \right) \hat{\Sigma}_\ell \right] \\
& + x^3 \left(-16\pi\hat{S}_\ell - \frac{31}{6}\pi\delta\hat{\Sigma}_\ell + \left[-\frac{5239}{504} - \frac{43}{2}\nu + \frac{41}{16}\delta\kappa_- + \kappa_+ \left(-\frac{271}{112} - \frac{43}{4}\nu \right) \right] \hat{S}_\ell^2 \right. \\
& + \left. \left\{ \delta \left[-\frac{817}{56} + \kappa_+ \left(-\frac{279}{56} - \frac{43}{4}\nu \right) - \frac{43}{2}\nu \right] + \kappa_- \left(\frac{279}{56} + \frac{1}{2}\nu \right) \right\} \hat{S}_\ell \hat{\Sigma}_\ell \right. \\
& + \left. \left[-\frac{25}{8} + \frac{344}{21}\nu + \frac{43}{2}\nu^2 + \delta\kappa_- \left(\frac{279}{112} + \frac{45}{16}\nu \right) + \kappa_+ \left(-\frac{279}{112} + \frac{243}{112}\nu + \frac{43}{4}\nu^2 \right) \right] \hat{\Sigma}_\ell^2 \right) \\
& + x^{7/2} \left\{ \left(\frac{476645}{6804} + \frac{6172}{189}\nu - \frac{2810}{27}\nu^2 \right) \hat{S}_\ell + \delta \left(\frac{9535}{336} + \frac{1849}{126}\nu - \frac{1501}{36}\nu^2 \right) \hat{\Sigma}_\ell + (16\pi + 8\pi\kappa_+) \hat{S}_\ell^2 \right. \\
& + \left. \left[-8\pi\kappa_- + \delta(16\pi + 8\pi\kappa_+) \right] \hat{S}_\ell \hat{\Sigma}_\ell + \left[\frac{1}{8}\pi - 4\pi\delta\kappa_- - 16\pi\nu + \kappa_+ (4\pi - 8\pi\nu) \right] \hat{\Sigma}_\ell^2 \right\} \\
& + x^4 \left(\left(-\frac{3485\pi}{96} + \frac{13879\pi}{72}\nu \right) \hat{S}_\ell + \delta \left(-\frac{7163\pi}{672} + \frac{130583\pi}{2016}\nu \right) \hat{\Sigma}_\ell + \left[-\frac{4289}{648} - \frac{295}{21}\nu + 54\nu^2 + \delta\kappa_- \left(\frac{935}{336} - \frac{2153}{144}\nu \right) \right. \right. \\
& + \left. \left. \kappa_+ \left(-\frac{124577}{9072} + \frac{3265}{126}\nu + 27\nu^2 \right) \right] \hat{S}_\ell^2 + \left\{ \kappa_- \left(\frac{74911}{4536} - \frac{52411}{1008}\nu + \frac{1181}{36}\nu^2 \right) \right. \right. \\
& + \left. \left. \delta \left[-\frac{160621}{9072} + \frac{9977}{252}\nu + 54\nu^2 + \kappa_+ \left(-\frac{74911}{4536} + \frac{41191}{1008}\nu + 27\nu^2 \right) \right] \right\} \hat{S}_\ell \hat{\Sigma}_\ell \right. \\
& + \left. \left[\frac{1633}{336} + \frac{465071}{18144}\nu - \frac{74033}{1008}\nu^2 - 54\nu^3 + \delta\kappa_- \left(\frac{74911}{9072} - \frac{46801}{2016}\nu + \frac{209}{144}\nu^2 \right) + \kappa_+ \left(-\frac{74911}{9072} + \frac{102979}{2592}\nu - \frac{7109}{168}\nu^2 - 27\nu^3 \right) \right] \hat{\Sigma}_\ell^2 \right) \Big].
\end{aligned} \tag{13}$$

As a consistency check, this result agrees with the GW flux computed for a (non-spinning) test body orbiting around a Kerr black hole in [78], to the given PN order.

We then combine these results to derive the evolution of the orbital frequency, from which we infer the change in the orbital phase, $\phi = \int \Omega(t) dt$, using the TaylorT2 approximant, e.g. [32], yielding

$$\begin{aligned}
\phi_{\text{spin}} = & -\frac{x^{-5/2}}{32\nu} \left[x^{3/2} \left(\frac{235}{6}\hat{S}_\ell + \frac{125}{8}\delta\hat{\Sigma}_\ell \right) + x^2 \left\{ (-50 - 25\kappa_+) \hat{S}_\ell^2 \right. \right. \\
& + \left. \left[25\kappa_- + \delta(-50 - 25\kappa_+) \right] \hat{S}_\ell \hat{\Sigma}_\ell + \left[-\frac{5}{16} + 50\nu + \frac{25}{2}\delta\kappa_- + \left(-\frac{25}{2} + 25\nu \right) \kappa_+ \right] \hat{\Sigma}_\ell^2 \right\} \\
& + x^{5/2} \log x \left[\left(-\frac{554345}{2016} - \frac{55}{8}\nu \right) \hat{S}_\ell + \left(-\frac{41745}{448} + \frac{15}{8}\nu \right) \delta\hat{\Sigma}_\ell \right] \\
& + x^3 \left(\frac{940}{3}\pi\hat{S}_\ell + \frac{745}{6}\pi\delta\hat{\Sigma}_\ell + \left[-\frac{31075}{126} + 60\nu + \frac{2215}{48}\delta\kappa_- + \left(\frac{15635}{84} + 30\nu \right) \kappa_+ \right] \hat{S}_\ell^2 \right. \\
& + \left. \left\{ \left(-\frac{47035}{336} - \frac{2575}{12}\nu \right) \kappa_- + \delta \left[-\frac{9775}{42} + 60\nu + \left(\frac{47035}{336} + 30\nu \right) \kappa_+ \right] \right\} \hat{S}_\ell \hat{\Sigma}_\ell \right. \\
& + \left. \left[-\frac{410825}{2688} + \frac{23535}{112}\nu - 60\nu^2 + \left(-\frac{47035}{672} - \frac{2935}{48}\nu \right) \delta\kappa_- + \left(\frac{47035}{672} - \frac{4415}{56}\nu - 30\nu^2 \right) \kappa_+ \right] \hat{\Sigma}_\ell^2 \right) \\
& + x^{7/2} \left\{ \left(-\frac{8980424995}{6096384} + \frac{6586595}{6048}\nu - \frac{305}{288}\nu^2 \right) \hat{S}_\ell + \left(-\frac{170978035}{387072} + \frac{2876425}{5376}\nu + \frac{4735}{1152}\nu^2 \right) \delta\hat{\Sigma}_\ell \right. \\
& + \left. \left(-100\pi - 50\pi\kappa_+ \right) \hat{S}_\ell^2 + \left[50\pi\kappa_- + \delta(-100\pi - 50\pi\kappa_+) \right] \hat{S}_\ell \hat{\Sigma}_\ell + \left[-\frac{15}{16}\pi + 100\nu\pi + 25\pi\delta\kappa_- + \left(-25\pi + 50\nu\pi \right) \kappa_+ \right] \hat{\Sigma}_\ell^2 \right\} \\
& + x^4 \left(\left(\frac{2388425\pi}{3024} - \frac{9925\pi}{36}\nu \right) \hat{S}_\ell + \delta \left(\frac{3237995\pi}{12096} - \frac{258245\pi}{2016}\nu \right) \hat{\Sigma}_\ell + \left[-\frac{83427805}{72576} - \frac{19720}{63}\nu + \frac{475}{24}\nu^2 + \left(\frac{3284125}{24192} + \frac{1115}{72}\nu \right) \delta\kappa_- \right. \right. \\
& + \left. \left. \left(\frac{55124675}{145152} - \frac{32825}{756}\nu + \frac{475}{48}\nu^2 \right) \kappa_+ \right] \hat{S}_\ell^2 + \left\{ \left(-\frac{35419925}{145152} - \frac{975955}{2016}\nu - \frac{10345}{144}\nu^2 \right) \kappa_- \right. \right.
\end{aligned} \tag{14}$$

$$\begin{aligned}
& + \delta \left[-\frac{66536845}{72576} - \frac{109535}{378} \nu + \frac{475}{24} \nu^2 + \left(\frac{35419925}{145152} - \frac{89065}{1512} \nu + \frac{475}{48} \nu^2 \right) \kappa_+ \right] \hat{S}_\ell \hat{\Sigma}_\ell \\
& + \left[-\frac{17815050265}{48771072} + \frac{26426305}{41472} \nu + \frac{12570535}{48384} \nu^2 - \frac{475}{24} \nu^3 + \left(-\frac{35419925}{290304} - \frac{2571605}{24192} \nu - \frac{5885}{288} \nu^2 \right) \delta \kappa_- \right. \\
& \left. + \left(\frac{35419925}{290304} - \frac{19990295}{145152} \nu + \frac{479845}{6048} \nu^2 - \frac{475}{48} \nu^3 \right) \kappa_+ \right] \hat{\Sigma}_\ell^2 \Big],
\end{aligned}$$

included in the ancillary file for the reader's convenience. We find perfect agreement in the overlap at linear order in the spin to 3.5PN of [58, 71], whereas spin-spin effects including finite-size corrections, both at 3.5PN and 4PN orders, are reported here for the first time.

Conclusions and Outlook. We have completed the knowledge of spin effects in the orbital phase evolution of compact binary systems to N²LO in the PN expansion of general relativity and quadratic order in the spins, corresponding to an overall 4PN order for rapidly rotating bodies, including both finite-size as well as tail effects. The various ingredients for the full derivation, such as

the gravitational potential and multipole moments, were obtained through the worldline EFT framework for spinning compact objects [31], which systematizes the two-body problem into a series of Feynman diagrams involving potential and radiation modes. Agreement is found in the overlap with previous PN results in the conservative [72–77] and radiation sectors [57, 58, 71]. In order to evaluate the impact of the new spin-dependent terms in the GW phase evolution, we used the leading quadrupolar approximation ($\phi_{\text{GW}} \simeq 2\phi$) to estimate the number of GW cycles in future detector's bands operating at design-sensitivity. The results, particularized to LISA (0.1mHz to $\min(f_{\text{ISCO}}, 1 \text{ Hz})$) and ET (1Hz to f_{ISCO} , with $f_{\text{ISCO}} = \frac{1}{6^{3/2} \pi G M}$), are summarized here:¹

LISA	$10^6 M_\odot + 10 M_\odot$	$10^5 M_\odot + 10 M_\odot$	$10^4 M_\odot + 10^4 M_\odot$
3.5PN	$224615. \hat{S}_\ell + (47903.6 + 23951.8 \kappa_+) \hat{S}_\ell^2 + 67352.3 \hat{\Sigma}_\ell + (47902.6 - 23951.8 \kappa_- + 23951.3 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (448.6 - 11975.6 \kappa_- + 11975.6 \kappa_+) \hat{\Sigma}_\ell^2$	$24002.2 \hat{S}_\ell + (5119.3 + 2559.6 \kappa_+) \hat{S}_\ell^2 + 7195.6 \hat{\Sigma}_\ell + (5118.2 - 2559.6 \kappa_- + 2559.1 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (47.5 - 1279.6 \kappa_- + 1279.6 \kappa_+) \hat{\Sigma}_\ell^2$	$4.7 \hat{S}_\ell + (1.2 + 0.6 \kappa_+) \hat{S}_\ell^2 - 0.6 \kappa_- \hat{S}_\ell \hat{\Sigma}_\ell + (-0.3 + 0.2 \kappa_+) \hat{\Sigma}_\ell^2$
4PN	$-164123. \hat{S}_\ell + (76034.4 - 8979. \kappa_- - 25119.6 \kappa_+) \hat{S}_\ell^2 - 55624.3 \hat{\Sigma}_\ell + (60639.1 + 16140.8 \kappa_- - 16140.1 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (24160.6 + 8070.1 \kappa_- - 8070.1 \kappa_+) \hat{\Sigma}_\ell^2$	$-16758.5 \hat{S}_\ell + (7764.3 - 916.7 \kappa_- - 2565. \kappa_+) \hat{S}_\ell^2 - 5678.7 \hat{\Sigma}_\ell + (6191.1 + 1648.5 \kappa_- - 1647.8 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (824. \kappa_- - 824. \kappa_+) \hat{\Sigma}_\ell^2$	$-2.8 \hat{S}_\ell + (1.5 - 0.4 \kappa_+) \hat{S}_\ell^2 + 0.4 \kappa_- \hat{S}_\ell \hat{\Sigma}_\ell + (0.2 - 0.1 \kappa_+) \hat{\Sigma}_\ell^2$
ET	$100 M_\odot + 1.4 M_\odot$	$10 M_\odot + 10 M_\odot$	$10 M_\odot + 1.4 M_\odot$
3.5PN	$163.2 \hat{S}_\ell + (35.2 + 17.6 \kappa_+) \hat{S}_\ell^2 + 47.3 \hat{\Sigma}_\ell + (34.2 - 17.6 \kappa_- + 17.1 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (-0.1 - 8.5 \kappa_- + 8.5 \kappa_+) \hat{\Sigma}_\ell^2$	$7.7 \hat{S}_\ell + (2. + \kappa_+) \hat{S}_\ell^2 - \kappa_- \hat{S}_\ell \hat{\Sigma}_\ell + (-0.5 + 0.2 \kappa_+) \hat{\Sigma}_\ell^2$	$20.5 \hat{S}_\ell + (4.7 + 2.4 \kappa_+) \hat{S}_\ell^2 + 4.4 \hat{\Sigma}_\ell + (3.6 - 2.4 \kappa_- + 1.8 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (-0.5 - 0.9 \kappa_- + 0.9 \kappa_+) \hat{\Sigma}_\ell^2$
4PN	$-119.9 \hat{S}_\ell + (56 - 6.4 \kappa_- - 18.4 \kappa_+) \hat{S}_\ell^2 - 39.4 \hat{\Sigma}_\ell + (43.5 + 12.2 \kappa_- - 11.5 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (17.3 + 5.8 \kappa_- - 5.8 \kappa_+) \hat{\Sigma}_\ell^2$	$-6.1 \hat{S}_\ell + (3.3 - \kappa_+) \hat{S}_\ell^2 + \kappa_- \hat{S}_\ell \hat{\Sigma}_\ell + (0.5 - 0.2 \kappa_+) \hat{\Sigma}_\ell^2$	$-15. \hat{S}_\ell + (7.4 - 0.6 \kappa_- - 2.4 \kappa_+) \hat{S}_\ell^2 - 3.8 \hat{\Sigma}_\ell + (4.5 + 1.9 \kappa_- - 1.1 \kappa_+) \hat{S}_\ell \hat{\Sigma}_\ell + (1.8 + 0.6 \kappa_- - 0.7 \kappa_+) \hat{\Sigma}_\ell^2$

Although generically relevant for comparable masses, spin effects become more important for binaries with unequal masses, which are also expected to exhibit larger spins [13]. Moreover, terms quadratic in the spins, depending on the inner structure of compact bodies and nearby surroundings through the κ_\pm couplings [17, 18], account for a large portion of the accumulated GW cycles. We therefore expect N²LO spin terms to play an important role in providing accurate waveform models and reliable physical interpretation of GW signals from ro-

tating compact binaries, thus motivating further in-depth studies to fully characterize their impact in detection and parameter estimation with forthcoming third-generation GW experiments.

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¹ We have not included the known absorption [79–81], radiation-reaction [82, 83] or cubic-in-spin [84] effects.

Supplemental Material

In what follows we summarize various results quoted in the main text. All these values are also conveniently given in the ancillary file. We use the (Euclidean) abbreviation

$$(ab) \equiv a^i b^i, \quad (a, b, c) \equiv \epsilon^{ijk} a^i b^j c^k, \quad (15)$$

to simplify the notation. We also define the norm vector as $\mathbf{n} \equiv \mathbf{r}/r$ and use $\mathbf{v} \equiv \mathbf{v}_1 - \mathbf{v}_2$ for the relative velocity. The coefficients of the spin-orbit and spin-spin binding energy in (3)-(4) are give by:

$$\begin{aligned}
e_3^0 &= -(n, S, v) - \delta(n, \Sigma, v), \\
e_5^0 &= 3\nu(nv)^2(n, S, v) + 3(1 + \nu)(n, S, v)v^2 + (-1 + 5\nu)\delta(n, \Sigma, v)v^2, \\
e_5^1 &= -4\nu(n, S, v) - 3\nu\delta(n, \Sigma, v), \\
e_7^0 &= 15\nu(-1 + 3\nu)(nv)^4(n, S, v) + 15\nu^2\delta(nv)^4(n, \Sigma, v) + 6(1 - 17\nu)\nu(nv)^2(n, S, v)v^2 \\
&\quad - 12\nu(2 + \nu)\delta(nv)^2(n, \Sigma, v)v^2 + (21 - 31\nu - 55\nu^2)(n, S, v)v^4 \\
&\quad + (-3 + 53\nu - 99\nu^2)\delta(n, \Sigma, v)v^4 \\
e_7^1 &= 2\nu(187 + 20\nu)(nv)^2(n, S, v) + \nu(179 + 12\nu)\delta(nv)^2(n, \Sigma, v) \\
&\quad + 2(24 - 143\nu + 40\nu^2)(n, S, v)v^2 + (-16 - 145\nu + 82\nu^2)\delta(n, \Sigma, v)v^2, \\
e_7^2 &= (6 - 111\nu - 8\nu^2)(n, S, v) - 2(-3 + 25\nu + 3\nu^2)\delta(n, \Sigma, v), \\
e_4^0 &= 6(2 + \kappa_+)(nS)^2 + [-6\kappa_- + 6\delta(2 + \kappa_+)](nS)(n\Sigma) \\
&\quad + [-3\delta\kappa_- + 3\kappa_+ - 6\nu(2 + \kappa_+)](n\Sigma)^2 - 2(2 + \kappa_+)\mathbf{S}^2 \\
&\quad + [2\kappa_- - 2\delta(2 + \kappa_+)](S\Sigma) + [\delta\kappa_- - \kappa_+ + 2\nu(2 + \kappa_+)]\Sigma^2, \\
e_6^0 &= (28 + 6\delta\kappa_- - 6\kappa_+)(vS)^2 + 4[(3 - 6\nu)\kappa_- + \delta(13 - 3\kappa_+)](vS)(v\Sigma) \\
&\quad + [-2\nu(38 + 3\delta\kappa_- - 9\kappa_+) + 6(4 + \delta\kappa_- - \kappa_+)](v\Sigma)^2 \\
&\quad + 6[-14 - 3\delta\kappa_- + 3\kappa_+ + 2\nu(2 + \kappa_+)](nv)(vS)(nS) + 6\{(-3 + 5\nu)\kappa_- \\
&\quad + \delta[-14 + 3\kappa_+ + \nu(2 + \kappa_+)]\}(nv)(v\Sigma)(nS) + \{-30\nu(2 + \kappa_+)(nv)^2 \\
&\quad - 6[-10 + 3\kappa_+ + \nu(2 + \kappa_+)]v^2\}(nS)^2 + 6[(-3 + 5\nu)\kappa_- \\
&\quad + 3\delta(-4 + \kappa_+) + \nu\delta(2 + \kappa_+)](nv)(vS)(n\Sigma) \\
&\quad - 6[-2\nu(19 + \delta\kappa_- - 4\kappa_+) + 3(4 + \delta\kappa_- - \kappa_+) + 2\nu^2(2 + \kappa_+)](nv)(v\Sigma)(n\Sigma) \\
&\quad + \{30\nu[\kappa_- - \delta(2 + \kappa_+)](nv)^2 - 6[-(3 + \nu)\kappa_- + 3\delta(-6 + \kappa_+) \\
&\quad + \nu\delta(2 + \kappa_+)]v^2\}(nS)(n\Sigma) + \{15\nu[\delta\kappa_- - \kappa_+ + 2\nu(2 + \kappa_+)](nv)^2 \\
&\quad + [48 + 9\delta\kappa_- - 9\kappa_+ + 6\nu^2(2 + \kappa_+) + 3\nu(-52 + \delta\kappa_- + 5\kappa_+)]v^2\}(n\Sigma)^2 \\
&\quad + \{6[4 + \delta\kappa_- - \kappa_+ + \nu(2 + \kappa_+)](nv)^2 + [-28 - 2\delta\kappa_- + 8\kappa_+ + 2\nu(2 + \kappa_+)]v^2\}\mathbf{S}^2 \\
&\quad + \{6\{(2 - 5\nu)\kappa_- + \delta[8 - 2\kappa_+ + \nu(2 + \kappa_+)]\}(nv)^2 + 2\{(-5 + 3\nu)\kappa_- \\
&\quad + \delta[-26 + 5\kappa_+ + \nu(2 + \kappa_+)]\}v^2\}(S\Sigma) + \{-3[-8 - 2\delta\kappa_- \\
&\quad + \nu(24 + 3\delta\kappa_- - 7\kappa_+) + 2\kappa_+ + 2\nu^2(2 + \kappa_+)](nv)^2 \\
&\quad + [-24 - 5\delta\kappa_- + \nu(76 + \delta\kappa_- - 11\kappa_+) + 5\kappa_+ - 2\nu^2(2 + \kappa_+)]v^2\}\Sigma^2, \\
e_6^1 &= (-36 + 9\delta\kappa_- - 15\kappa_+)(nS)^2 - 4(8\delta - 6\kappa_- + 9\nu\kappa_- + 6\delta\kappa_+)(nS)(n\Sigma) \\
&\quad + [12(\delta\kappa_- - \kappa_+) + \nu(30 - 9\delta\kappa_- + 33\kappa_+)](n\Sigma)^2 + (8 - 3\delta\kappa_- + 5\kappa_+)\mathbf{S}^2 \\
&\quad + [4(-2 + 3\nu)\kappa_- + 8\delta(1 + \kappa_+)](S\Sigma) \\
&\quad + [-4\delta\kappa_- + \nu(-10 + 3\delta\kappa_- - 11\kappa_+) + 4\kappa_+]\Sigma^2, \\
e_8^0 &= \{-6[-8 - 3\delta\kappa_- + 3\kappa_+ + 6\nu^2(2 + \kappa_+) - 2\nu(-26 + 6\delta\kappa_- + 7\kappa_+)](nv)^2 \\
&\quad - 4[-39 - 9\delta\kappa_- + \nu(70 + 27\delta\kappa_- - 15\kappa_+) + 9\kappa_+]v^2\}(vS)^2 \\
&\quad + \{6[8 + 3\delta\kappa_- - 2\nu(12 + 2\delta\kappa_- - 5\kappa_+) - 3\kappa_+ \\
&\quad + 6\nu^3(2 + \kappa_+) + \nu^2(-44 - 9\delta\kappa_- + 7\kappa_+)](nv)^2 + 4[\nu^2(286 + 27\delta\kappa_- - 69\kappa_+) \\
&\quad - 3\nu(65 + 10\delta\kappa_- - 16\kappa_+) + 9(4 + \delta\kappa_- - \kappa_+)]v^2\}(v\Sigma)^2 \\
&\quad + \{60\nu[10 - 3\delta\kappa_- - 5\kappa_+ + 6\nu(2 + \kappa_+)](nv)^3 \\
&\quad - 6[86 + 21\delta\kappa_- - 2\nu(90 + 30\delta\kappa_- - 11\kappa_+) - 21\kappa_+ + 54\nu^2(2 + \kappa_+)](nv)v^2\}(vS)(nS) \\
&\quad + \{60\nu[\kappa_- + 3\nu\kappa_- + (-1 + 3\nu)\delta(2 + \kappa_+)](nv)^3
\end{aligned}$$

$$\begin{aligned}
& -6 \{ (21 - 83\nu + 93\nu^2) \kappa_- + \delta [86 - 21\kappa_+ + 27\nu^2 (2 + \kappa_+)] \\
& + \nu (-214 + 41\kappa_+) \} (nv) v^2 (v\Sigma) (nS) + \{ -210\nu (-1 + 3\nu) (2 + \kappa_+) (nv)^4 \\
& + 60\nu [-18 - \kappa_+ + 15\nu (2 + \kappa_+)] (nv)^2 v^2 + 6 [54 - 21\kappa_+ + 27\nu^2 (2 + \kappa_+)] \\
& + \nu (-114 + 23\kappa_+) v^4 \} (nS)^2 + (60\nu \{ \kappa_- - \delta\kappa_+ + 3\nu [\kappa_- \\
& + \delta(2 + \kappa_+)] \} (nv)^3 - 6 \{ (21 - 83\nu + 93\nu^2) \kappa_- + \delta [80 - 21\kappa_+ + 27\nu^2 (2 + \kappa_+)] \\
& + \nu (-174 + 41\kappa_+) \} (nv) v^2 (vS) (n\Sigma) + \{ -60\nu [-\delta\kappa_- + \kappa_+ + 6\nu^2 (2 + \kappa_+)] \\
& - 2\nu (7 + \kappa_+) \} (nv)^3 + 6 [-80 - 21\delta\kappa_- + 2\nu (207 + 31\delta\kappa_- - 52\kappa_+) + 21\kappa_+ \\
& + 54\nu^3 (2 + \kappa_+) + \nu^2 (-596 - 33\delta\kappa_- + 115\kappa_+)] (nv) v^2 \} (v\Sigma) (n\Sigma) \\
& + (210\nu (-1 + 3\nu) [\kappa_- - \delta(2 + \kappa_+)] (nv)^4 + 60\nu \{ \kappa_- - 15\nu\kappa_- \\
& + \delta [-10 - \kappa_+ + 15\nu (2 + \kappa_+)] \} (nv)^2 v^2 \\
& + 6 \{ (21 - 23\nu - 27\nu^2) \kappa_- + \delta [102 - 21\kappa_+ + 27\nu^2 (2 + \kappa_+)] \\
& + \nu (-258 + 23\kappa_+) \} v^4) (nS) (n\Sigma) + \{ 105\nu (-1 + 3\nu) [\delta\kappa_- - \kappa_+ + 2\nu (2 + \kappa_+)] (nv)^4 \\
& - 30\nu (-1 + 15\nu) [\delta\kappa_- - \kappa_+ + 2\nu (2 + \kappa_+)] (nv)^2 v^2 \\
& - 3 [\nu (524 + 23\delta\kappa_- - 65\kappa_+) - 3 (32 + 7\delta\kappa_- - 7\kappa_+) + 54\nu^3 (2 + \kappa_+)] \\
& + \nu^2 (-804 + 27\delta\kappa_- + 19\kappa_+) v^4 \} (n\Sigma)^2 + \{ 30\nu [-10 + 2\delta\kappa_- + \kappa_+ + 3\nu (2 + \kappa_+)] (nv)^4 \\
& - 12 [\nu (-10 + 12\delta\kappa_- - 3\kappa_+) + 15\nu^2 (2 + \kappa_+) + 3 (-4 - \delta\kappa_- + \kappa_+)] (nv)^2 v^2 \\
& - 2 [78 + 6\delta\kappa_- - 27\kappa_+ + 27\nu^2 (2 + \kappa_+) + \nu (-154 - 18\delta\kappa_- + 33\kappa_+)] v^4 \} \mathbf{S}^2 \\
& + (30\nu [\kappa_- - 11\nu\kappa_- + (-1 + 3\nu)\delta(2 + \kappa_+)] (nv)^4 + \{ 36(2 - 9\nu + 21\nu^2) \kappa_- \\
& - 12\delta [\nu (46 - 15\kappa_+) + 6(-4 + \kappa_+) + 15\nu^2 (2 + \kappa_+)] \} (nv)^2 v^2 \\
& - 2 \{ 3(11 - 25\nu + 15\nu^2) \kappa_- + \delta [150 - 33\kappa_+ + 27\nu^2 (2 + \kappa_+)] \\
& + \nu (-370 + 51\kappa_+) \} v^4) (S\Sigma) + \{ -15\nu [-\delta\kappa_- + \nu(12 + 7\delta\kappa_- - 9\kappa_+) \\
& + \kappa_+ + 6\nu^2 (2 + \kappa_+)] (nv)^4 + 6 [3\nu^2 (68 + 13\delta\kappa_- - 23\kappa_+) + 6(4 + \delta\kappa_- - \kappa_+) \\
& + 30\nu^3 (2 + \kappa_+) + \nu (-128 - 21\delta\kappa_- + 33\kappa_+)] (nv)^2 v^2 \\
& + [-144 - 33\delta\kappa_- + 3\nu(260 + 21\delta\kappa_- - 43\kappa_+) + 33\kappa_+ + 54\nu^3 (2 + \kappa_+)] \\
& + \nu^2 (-1172 - 9\delta\kappa_- + 111\kappa_+) v^4 \} \mathbf{\Sigma}^2 \\
& + (vS) \left[(-12 \{ (-3 + 7\nu + 21\nu^2) \kappa_- + \delta [-8 - \nu(-2 + \kappa_+) \right. \\
& + 3\kappa_+ + 3\nu^2 (2 + \kappa_+)] \} (nv)^2 + 4 \{ 6(3 - 13\nu + 18\nu^2) \kappa_- \\
& + \delta [75 - 18\kappa_+ + 2\nu(-89 + 21\kappa_+)] \} v^2) (v\Sigma) \Big], \\
e_8^1 &= 4 [\nu (-4 + 3\delta\kappa_- - 41\kappa_+) + 5(2 + \delta\kappa_- + \kappa_+)] (vS)^2 \\
& + 8 \{ -6(-2 + \nu)\nu\kappa_- + \delta [5 + \nu(6 - 22\kappa_+)] \} (vS) (v\Sigma) \\
& + 4\nu [20 + 17\delta\kappa_- - 17\kappa_+ + \nu(-38 - 3\delta\kappa_- + 47\kappa_+)] (v\Sigma)^2 \\
& + [72\nu^2 (2 + \kappa_+) - 8(5 + 21\delta\kappa_- + 21\kappa_+) + 2\nu(930 + 15\delta\kappa_- + 649\kappa_+)] (nv) (vS) (nS) \\
& + \{ -2\nu(149 + 48\nu)\kappa_- + \delta [-56 + 36\nu^2 (2 + \kappa_+) + \nu(824 + 634\kappa_+)] \} (nv) (v\Sigma) (nS) \\
& + \{ -6 [26 - 43\delta\kappa_- - 43\kappa_+ + 40\nu^2 (2 + \kappa_+) + \nu(740 + 27\delta\kappa_- + 340\kappa_+)] (nv)^2 \\
& - 2 [-54 + 46\delta\kappa_- + \nu(-682 + 9\delta\kappa_- - 348\kappa_+) + 46\kappa_+ + 12\nu^2 (2 + \kappa_+)] v^2 \} (nS)^2 \\
& + \{ -2\nu(149 + 48\nu)\kappa_- + \delta [-32 + 36\nu^2 (2 + \kappa_+) + \nu(876 + 634\kappa_+)] \} (nv) (vS) (n\Sigma) \\
& - 2\nu [84 + 233\delta\kappa_- - 233\kappa_+ + 36\nu^2 (2 + \kappa_+) + \nu(730 + 33\delta\kappa_- + 601\kappa_+)] (nv) (v\Sigma) (n\Sigma) \\
& + (6\nu(141 + 148\nu)\kappa_- - 6\nu\delta [752 + 313\kappa_+ + 40\nu(2 + \kappa_+)] \} (nv)^2 \\
& + \{ 2\nu(-173 + 48\nu)\kappa_- + \delta [64 - 24\nu^2 (2 + \kappa_+) + 6\nu(246 + 119\kappa_+)] \} v^2) (nS) (n\Sigma) \\
& + \{ 3\nu [-60 + 227\delta\kappa_- - 227\kappa_+ + 80\nu^2 (2 + \kappa_+) + 2\nu(768 + 47\delta\kappa_- + 266\kappa_+)] (nv)^2 \\
& + \nu [156 - 265\delta\kappa_- + 2\nu(-824 + 15\delta\kappa_- - 372\kappa_+) + 265\kappa_+ + 24\nu^2 (2 + \kappa_+)] v^2 \} (n\Sigma)^2 \\
& + \{ 2 [28\nu^2 (2 + \kappa_+) - 5(-2 + 3\delta\kappa_- + 3\kappa_+) + \nu(658 + 22\delta\kappa_- + 209\kappa_+)] (nv)^2 \\
& + [2\nu(-302 + \delta\kappa_- - 110\kappa_+) + 8\nu^2 (2 + \kappa_+) + 4(-11 + 6\delta\kappa_- + 6\kappa_+)] v^2 \} \mathbf{S}^2 \\
& + (\{ -2\nu(127 + 116\nu)\kappa_- + 2\delta [16 + 28\nu^2 (2 + \kappa_+) + \nu(702 + 187\kappa_+)] \} (nv)^2 \\
& + 2 \{ (63 - 8\nu)\nu\kappa_- + \delta [-24 + 4\nu^2 (2 + \kappa_+) - \nu(338 + 111\kappa_+)] \} v^2) (S\Sigma) \\
& + \{ -\nu [-124 + 157\delta\kappa_- - 157\kappa_+ + 56\nu^2 (2 + \kappa_+) + 2\nu(778 + 36\delta\kappa_- + 151\kappa_+)] (nv)^2 \\
& - \nu [108 - 87\delta\kappa_- + \nu(-808 + 6\delta\kappa_- - 228\kappa_+) + 87\kappa_+ + 8\nu^2 (2 + \kappa_+)] v^2 \} \mathbf{\Sigma}^2,
\end{aligned}$$

$$\begin{aligned}
e_8^2 = & [72 - 414 \delta \kappa_- - 120 \kappa_+ + 94 \nu (50 + 17 \kappa_+)] (nS)^2 \\
& + 2 \{ (-147 + 29 \nu) \kappa_- + \delta [-70 + 147 \kappa_+ + \nu (2322 + 799 \kappa_+)] \} (nS) (n\Sigma) \\
& + [147 (-\delta \kappa_- + \kappa_+) + \nu (420 - 385 \delta \kappa_- + 91 \kappa_+) - 2 \nu^2 (2280 + 799 \kappa_+)] (n\Sigma)^2 \\
& + 2 [72 + 69 \delta \kappa_- + 20 \kappa_+ - \nu (662 + 341 \kappa_+)] \mathbf{S}^2 \\
& + \{ 2 (49 + 65 \nu) \kappa_- - 2 \delta [-70 + 49 \kappa_+ + \nu (634 + 341 \kappa_+)] \} (S\Sigma) \\
& + [7 \nu (4 + 29 \delta \kappa_- - 15 \kappa_+) + 49 (\delta \kappa_- - \kappa_+) + 2 \nu^2 (592 + 341 \kappa_+)] \Sigma^2 .
\end{aligned}$$

On the other hand, for the radiated flux in (8)-(9) we find

$$\begin{aligned}
f_3^0 = & -78 (nv)^2 (n, S, v) - 51 \delta (nv)^2 (n, \Sigma, v) + 80 (n, S, v) v^2 + 43 \delta (n, \Sigma, v) v^2 , \\
f_3^1 = & 2 (n, S, v) - \delta (n, \Sigma, v) , \\
f_5^0 = & 48 (187 - 262 \nu) (nv)^4 (n, S, v) + 3 (2647 - 3568 \nu) \delta (nv)^4 (n, \Sigma, v) \\
& + 36 (-391 + 556 \nu) (nv)^2 (n, S, v) v^2 + 12 (-788 + 1207 \nu) \delta (nv)^2 (n, \Sigma, v) v^2 \\
& + (4828 - 7240 \nu) (n, S, v) v^4 + (2603 - 4160 \nu) \delta (n, \Sigma, v) v^4 \\
f_5^1 = & 22 (293 - 46 \nu) (nv)^2 (n, S, v) + (7327 - 2398 \nu) \delta (nv)^2 (n, \Sigma, v) \\
& + 2 (-3805 + 224 \nu) (n, S, v) v^2 + (-5387 + 994 \nu) \delta (n, \Sigma, v) v^2 \\
f_5^2 = & -2 (472 + 195 \nu) (n, S, v) + (137 - 238 \nu) \delta (n, \Sigma, v) , \\
f_7^0 = & -10 (2939 - 26896 \nu + 6632 \nu^2) (nv)^6 (n, S, v) \\
& + 5 (12215 + 56391 \nu - 23312 \nu^2) \delta (nv)^6 (n, \Sigma, v) \\
& + 6 (11731 - 97552 \nu + 41208 \nu^2) (nv)^4 (n, S, v) v^2 \\
& + 3 (-28423 - 177451 \nu + 98604 \nu^2) \delta (nv)^4 (n, \Sigma, v) v^2 \\
& - 6 (10485 - 67206 \nu + 44150 \nu^2) (nv)^2 (n, S, v) v^4 \\
& + 3 (4369 + 99655 \nu - 75750 \nu^2) \delta (nv)^2 (n, \Sigma, v) v^4 \\
& + 2 (11237 - 43924 \nu + 41766 \nu^2) (n, S, v) v^6 + (8847 - 49916 \nu + 50198 \nu^2) \delta (n, \Sigma, v) v^6 , \\
f_7^1 = & -2 (36959 - 204812 \nu + 94378 \nu^2) (nv)^4 (n, S, v) \\
& + (-140184 + 619741 \nu - 243534 \nu^2) \delta (nv)^4 (n, \Sigma, v) \\
& + 2 (87510 - 279467 \nu + 111074 \nu^2) (nv)^2 (n, S, v) v^2 \\
& + (177222 - 664927 \nu + 273594 \nu^2) \delta (nv)^2 (n, \Sigma, v) v^2 \\
& - 2 (44009 - 72691 \nu + 23602 \nu^2) (n, S, v) v^4 - 2 (33267 - 65918 \nu + 26015 \nu^2) \delta (n, \Sigma, v) v^4 , \\
f_7^2 = & (165714 - 210143 \nu - 82280 \nu^2) (nv)^2 (n, S, v) \\
& + (-169603 + 142753 \nu - 108405 \nu^2) \delta (nv)^2 (n, \Sigma, v) \\
& + (98244 + 77017 \nu + 23098 \nu^2) (n, S, v) v^2 + 2 (76462 - 19183 \nu + 14961 \nu^2) \delta (n, \Sigma, v) v^2 , \\
f_7^3 = & 3 (67927 + 32320 \nu - 2362 \nu^2) (n, S, v) + (21076 + 28209 \nu - 6021 \nu^2) \delta (n, \Sigma, v) \\
f_4^0 = & 72 (2 + \kappa_+) (vS)^2 + 72 [-\kappa_- + \delta (2 + \kappa_+)] (vS) (v\Sigma) \\
& + [1 - 36 \delta \kappa_- + 36 \kappa_+ - 72 \nu (2 + \kappa_+)] (v\Sigma)^2 - 348 (2 + \kappa_+) (nv) (vS) (nS) \\
& + 174 [\kappa_- - \delta (2 + \kappa_+)] (nv) (v\Sigma) (nS) + [816 (2 + \kappa_+) (nv)^2 - 504 (2 + \kappa_+) v^2] (nS)^2 \\
& + 174 [\kappa_- - \delta (2 + \kappa_+)] (nv) (vS) (n\Sigma) \\
& + 6 [-1 + 29 \delta \kappa_- - 29 \kappa_+ + 58 \nu (2 + \kappa_+)] (nv) (v\Sigma) (n\Sigma) \\
& + \{ 816 [-\kappa_- + \delta (2 + \kappa_+)] (nv)^2 + 504 [\kappa_- - \delta (2 + \kappa_+)] v^2 \} (nS) (n\Sigma) \\
& + \{ [9 - 408 \delta \kappa_- + 408 \kappa_+ - 816 \nu (2 + \kappa_+)] (nv)^2 + 252 [\delta \kappa_- - \kappa_+ + 2 \nu (2 + \kappa_+)] v^2 \} (n\Sigma)^2 \\
& + [-156 (2 + \kappa_+) (nv)^2 + 144 (2 + \kappa_+) v^2] \mathbf{S}^2 + \{ 156 [\kappa_- - \delta (2 + \kappa_+)] (nv)^2 \\
& + 144 [-\kappa_- + \delta (2 + \kappa_+)] v^2 \} (S\Sigma) + \{ [9 + 78 \delta \kappa_- - 78 \kappa_+ + 156 \nu (2 + \kappa_+)] (nv)^2 \\
& + [3 - 72 \delta \kappa_- + 72 \kappa_+ - 144 \nu (2 + \kappa_+)] v^2 \} \Sigma^2 , \\
f_6^0 = & \{ [3897 + 3189 \delta \kappa_- - 4557 \kappa_+ + 2550 \nu (2 + \kappa_+)] (nv)^2 \\
& + [-3239 - 834 \delta \kappa_- - 162 \kappa_+ + 60 \nu (2 + \kappa_+)] v^2 \} (vS)^2 \\
& + \{ 6 \{ (1291 - 2551 \nu) \kappa_- + \delta [1921 - 1291 \kappa_+ + 425 \nu (2 + \kappa_+)] \} (nv)^2 \\
& + \{ 84 (-8 + 39 \nu) \kappa_- + 2 \delta [-2047 + 336 \kappa_+ + 30 \nu (2 + \kappa_+)] \} v^2 \} (vS) (v\Sigma) \\
& + \{ [7062 + 3873 \delta \kappa_- - 6 \nu (3036 + 744 \delta \kappa_- - 2035 \kappa_+) - 3873 \kappa_+ - 2550 \nu^2 (2 + \kappa_+)] (nv)^2 \\
& + [-1121 - 336 \delta \kappa_- + 2 \nu (2447 + 402 \delta \kappa_- - 738 \kappa_+) + 336 \kappa_+ - 60 \nu^2 (2 + \kappa_+)] v^2 \} (v\Sigma)^2 \\
& + \{ -12 [-686 + 823 \delta \kappa_- - 2159 \kappa_+ + 1216 \nu (2 + \kappa_+)] (nv)^3 \\
& + [-90 + 4761 \delta \kappa_- - 9945 \kappa_+ + 2046 \nu (2 + \kappa_+)] (nv) v^2 \} (vS) (nS)
\end{aligned}$$

$$\begin{aligned}
& + \left(-12 \{7(213 - 322\nu)\kappa_- + \delta[10 - 1491\kappa_+ + 608\nu(2 + \kappa_+)]\} (nv)^3 \right. \\
& + \left. \{57(129 - 185\nu)\kappa_- + 3\delta[-854 - 2451\kappa_+ + 341\nu(2 + \kappa_+)]\} (nv)v^2 \right) (vS)(nS) \\
& + \{15[-2719 + 431\delta\kappa_- - 3019\kappa_+ + 2384\nu(2 + \kappa_+)](nv)^4 \\
& - 6[-5435 + 677\delta\kappa_- - 6803\kappa_+ + 5248\nu(2 + \kappa_+)](nv)^2v^2 \\
& + [-3273 + 285\delta\kappa_- - 7389\kappa_+ + 6816\nu(2 + \kappa_+)]v^4\} (nS)^2 \\
& + \left(-3 \{28(213 - 322\nu)\kappa_- + \delta[2567 - 5964\kappa_+ + 2432\nu(2 + \kappa_+)]\} (nv)^3 \right. \\
& + \left. \{57(129 - 185\nu)\kappa_- + 3\delta[1365 - 2451\kappa_+ + 341\nu(2 + \kappa_+)]\} (nv)v^2 \right) (vS)(n\Sigma) \\
& + \{3[-5589 - 5964\delta\kappa_- + \nu(8143 + 5724\delta\kappa_- - 17652\kappa_+) \\
& + 5964\kappa_+ + 4864\nu^2(2 + \kappa_+)](nv)^3 \\
& - 3[-1297 - 2451\delta\kappa_- + \nu(1299 + 1928\delta\kappa_- - 6830\kappa_+) \\
& + 2451\kappa_+ + 682\nu^2(2 + \kappa_+)](nv)v^2\} (vS)(n\Sigma) \\
& + \left(30[(1725 - 2054\nu)\kappa_- + 1192\nu\delta(2 + \kappa_+) - 3\delta(394 + 575\kappa_+)](nv)^4 \right. \\
& - 24[17(110 - 117\nu)\kappa_- + 1312\nu\delta(2 + \kappa_+) - 374\delta(4 + 5\kappa_+)](nv)^2v^2 \\
& + 6\{(1279 - 1326\nu)\kappa_- + \delta[-722 - 1279\kappa_+ + 1136\nu(2 + \kappa_+)]\}v^4\} (nS)(n\Sigma) \\
& + \{-15[-614 - 1725\delta\kappa_- + \nu(-2032 + 1623\delta\kappa_- \\
& - 5073\kappa_+) + 1725\kappa_+ + 2384\nu^2(2 + \kappa_+)](nv)^4 \\
& + 6[\nu(-6634 + 3301\delta\kappa_- - 10781\kappa_+) - 4(169 + 935\delta\kappa_- \\
& - 935\kappa_+) + 5248\nu^2(2 + \kappa_+)](nv)^2v^2 \\
& + 3[515 + 1279\delta\kappa_- - 1279\kappa_+ - 2272\nu^2(2 + \kappa_+) \\
& + \nu(1814 - 1231\delta\kappa_- + 3789\kappa_+)]v^4\} (n\Sigma)^2 \\
& + \{3[5906 + 379\delta\kappa_- + 2153\kappa_+ - 2352\nu(2 + \kappa_+)](nv)^4 \\
& + 12[-2166 - 108\delta\kappa_- - 731\kappa_+ + 747\nu(2 + \kappa_+)](nv)^2v^2 \\
& + 3[3046 + 61\delta\kappa_- + 839\kappa_+ - 764\nu(2 + \kappa_+)]v^4\} \mathbf{S}^2 \\
& + \left(-6 \{887 - 418\nu\}\kappa_- + \delta[-4221 - 887\kappa_+ + 1176\nu(2 + \kappa_+)] \right) (nv)^4 \\
& + 12 \{7(89 - 45\nu)\kappa_- + \delta[-2903 - 623\kappa_+ + 747\nu(2 + \kappa_+)]\} (nv)^2v^2 \\
& + \{6(-389 + 260\nu)\kappa_- + \delta[10406 + 2334\kappa_+ - 2292\nu(2 + \kappa_+)]\}v^4\} (S\Sigma) \\
& + \{3[1619 - 887\delta\kappa_- + \nu(-10798 + 797\delta\kappa_- \\
& - 2571\kappa_+) + 887\kappa_+ + 2352\nu^2(2 + \kappa_+)](nv)^4 \\
& - 3[2381 - 1246\delta\kappa_- + 2\nu(-7193 + 531\delta\kappa_- - 1777\kappa_+) + 1246\kappa_+ + 2988\nu^2(2 + \kappa_+)](nv)^2v^2 \\
& + [1898 - 1167\delta\kappa_- + \nu(-11668 + 963\delta\kappa_- - 3297\kappa_+) \\
& + 1167\kappa_+ + 2292\nu^2(2 + \kappa_+)]v^4\} \mathbf{\Sigma}^2, \\
f_6^1 = & -2[6810 - 1166\delta\kappa_- + 3663\kappa_+ + 42\nu(2 + \kappa_+)](vS)^2 \\
& - 2\{(-4829 + 4622\nu)\kappa_- + \delta[4620 + 4829\kappa_+ + 42\nu(2 + \kappa_+)]\}(vS)(v\Sigma) \\
& + [2392 + 4829\delta\kappa_- - 4829\kappa_+ + 84\nu^2(2 + \kappa_+) + \nu(4904 - 2290\delta\kappa_- + 11948\kappa_+)](v\Sigma)^2 \\
& + 6[10250 - 1706\delta\kappa_- + 5338\kappa_+ + 161\nu(2 + \kappa_+)](nv)(vS)(nS) \\
& + [9(-2348 + 2221\nu)\kappa_- + 483\nu\delta(2 + \kappa_+) + 4\delta(6343 + 5283\kappa_+)](nv)(v\Sigma)(nS) \\
& + \{6[-25044 + 2399\delta\kappa_- - 11299\kappa_+ + 169\nu(2 + \kappa_+)](nv)^2 \\
& + [91376 - 6190\delta\kappa_- + 38868\kappa_+ - 996\nu(2 + \kappa_+)]v^2\}(nS)^2 \\
& + [9(-2348 + 2221\nu)\kappa_- + 483\nu\delta(2 + \kappa_+) + 4\delta(5804 + 5283\kappa_+)](nv)(vS)(n\Sigma) \\
& + [\nu(-34902 + 9753\delta\kappa_- - 52017\kappa_+) - 4(1226 + 5283\delta\kappa_- \\
& - 5283\kappa_+) - 966\nu^2(2 + \kappa_+)](nv)(v\Sigma)(n\Sigma) \\
& + \left(6 \{9(1522 - 1085\nu)\kappa_- + \delta[-23537 - 13698\kappa_+ + 169\nu(2 + \kappa_+)]\} (nv)^2 \right. \\
& + \left. \{2(-22529 + 12878\nu)\kappa_- + \delta[91858 + 45058\kappa_+ - 996\nu(2 + \kappa_+)]\} v^2 \right) (nS)(n\Sigma) \\
& + \{-3[934 - 13698\delta\kappa_- + \nu(-44112 + 4967\delta\kappa_- \\
& - 32363\kappa_+) + 13698\kappa_+ + 338\nu^2(2 + \kappa_+)](nv)^2 \\
& + [4850 - 22529\delta\kappa_- + \nu(-92766 + 6688\delta\kappa_- - 51746\kappa_+) \\
& + 22529\kappa_+ + 996\nu^2(2 + \kappa_+)]v^2\}(n\Sigma)^2 \\
& + \{[24436 - 1386\delta\kappa_- + 11922\kappa_+ - 660\nu(2 + \kappa_+)](nv)^2 \\
& + 2[643\delta\kappa_- + 180\nu(2 + \kappa_+) - 7(1494 + 751\kappa_+)]v^2\} \mathbf{S}^2
\end{aligned}$$

$$\begin{aligned}
& + \left(-12 \{ (1109 - 517\nu) \kappa_- + \delta [-2299 - 1109 \kappa_+ + 55\nu (2 + \kappa_+)] \} (nv)^2 \right. \\
& + \left. \{ 8 (1475 - 688\nu) \kappa_- + 4\delta [-6109 - 2950 \kappa_+ + 90\nu (2 + \kappa_+)] \} v^2 \right) (S\Sigma) \\
& + \{ [604 - 6654\delta \kappa_- + 4\nu (-7717 + 429\delta \kappa_- - 3756 \kappa_+) + 6654 \kappa_+ + 660\nu^2 (2 + \kappa_+)] (nv)^2 \\
& - 2 [1102 - 2950\delta \kappa_- + \nu (-14072 + 733\delta \kappa_- - 6633 \kappa_+) \\
& + 2950 \kappa_+ + 180\nu^2 (2 + \kappa_+)] v^2 \} \Sigma^2 \\
f_6^2 = & 3 [-39 - 2\delta \kappa_- - 18 \kappa_+ + 72\nu (2 + \kappa_+)] (nS)^2 + \{ 48 (1 - 4\nu) \kappa_- \\
& + \delta [-98 - 48 \kappa_+ + 216\nu (2 + \kappa_+)] \} (nS) (n\Sigma) - 2 [1 - 12\delta \kappa_- + \nu (-58 + 51\delta \kappa_- - 75 \kappa_+) \\
& + 12 \kappa_+ + 108\nu^2 (2 + \kappa_+)] (n\Sigma)^2 + 2 [23 + \delta \kappa_- + 9 \kappa_+ - 36\nu (2 + \kappa_+)] S^2 \\
& - 4 \{ 4 (1 - 4\nu) \kappa_- + \delta [-7 - 4 \kappa_+ + 18\nu (2 + \kappa_+)] \} (S\Sigma) \\
& + [3 - 8\delta \kappa_- + \nu (-48 + 34\delta \kappa_- - 50 \kappa_+) + 8 \kappa_+ + 72\nu^2 (2 + \kappa_+)] \Sigma^2 \\
f_8^0 = & \{ -3 [-88255 + 204790\delta \kappa_- + 21058 \kappa_+ + 14048\nu^2 (2 + \kappa_+) \\
& + \nu (47364 - 135196\delta \kappa_- + 25148 \kappa_+)] (nv)^4 + 3 [-52309 + 163221\delta \kappa_- + 39339 \kappa_+ \\
& + 24630\nu^2 (2 + \kappa_+) - 2\nu (-8060 + 59661\delta \kappa_- + 19254 \kappa_+)] (nv)^2 v^2 \\
& - 6 [12058 + 7347\delta \kappa_- + 3273 \kappa_+ + 4872\nu^2 (2 + \kappa_+) \\
& - 6\nu (1883 + 1279\delta \kappa_- + 1941 \kappa_+)] v^4 \} (vS)^2 \\
& + \left(-6 \{ 2 (45933 - 244876\nu + 131684\nu^2) \kappa_- + \delta [90509 - 91866 \kappa_+ + 7024\nu^2 (2 + \kappa_+) \right. \\
& + \left. \nu (-153137 + 80172 \kappa_+) \} \right) (nv)^4 \\
& + 6 \{ 3 (20647 - 122283\nu + 75443\nu^2) \kappa_- + \delta [70079 - 61941 \kappa_+ + 12315\nu^2 (2 + \kappa_+) \\
& + 3\nu (-48145 + 13469 \kappa_+)] \} (nv)^2 v^2 - 12 \{ 3 (679 - 4236\nu + 4304\nu^2) \kappa_- \\
& + \delta [10347 - 2037 \kappa_+ + 2436\nu^2 (2 + \kappa_+) - 2\nu (5947 + 993 \kappa_+)] \} v^4 \} (vS) (v\Sigma) \\
& + \{ 3 [-175013 - 91866\delta \kappa_- + \nu (581226 + 284962\delta \kappa_- - 468694 \kappa_+) \\
& + 91866 \kappa_+ + 14048\nu^3 (2 + \kappa_+) + \nu^2 (-638641 - 128172\delta \kappa_- + 288516 \kappa_+)] (nv)^4 \\
& + [9\nu^2 (189985 + 35669\delta \kappa_- - 62607 \kappa_+) - 18\nu (72279 + 33938\delta \kappa_- - 54585 \kappa_+) \\
& + 33 (10717 + 5631\delta \kappa_- - 5631 \kappa_+) - 73890\nu^3 (2 + \kappa_+)] (nv)^2 v^2 \\
& + 6 [\nu (31842 + 5361\delta \kappa_- - 9435 \kappa_+) - 3 (1934 + 679\delta \kappa_- - 679 \kappa_+) \\
& - 2\nu^2 (17741 + 2619\delta \kappa_- - 633 \kappa_+) + 4872\nu^3 (2 + \kappa_+)] v^4 \} (v\Sigma)^2 \\
& + \left(90 [-20466 + 22857\delta \kappa_- + 3619 \kappa_+ + 1232\nu^2 (2 + \kappa_+) + \nu (35394 - 14094\delta \kappa_- \right. \\
& + 8402 \kappa_+)] (nv)^5 - 12 [-176174 + 171751\delta \kappa_- + 50677 \kappa_+ + 32588\nu^2 (2 + \kappa_+) \\
& + \nu (333866 - 116344\delta \kappa_- + 22292 \kappa_+)] (nv)^3 v^2 + 3 \{ -132202 + 118101\delta \kappa_- \\
& + 46851 \kappa_+ + 89034\nu^2 (2 + \kappa_+) - 2\nu [54081\delta \kappa_- + 26 (-7777 + 492 \kappa_+)] \} (nv) v^4 \right) (vS) (nS) \\
& + \left(90 \{ (9619 - 56962\nu + 27572\nu^2) \kappa_- + \delta [1766 - 9619 \kappa_+ + 616\nu^2 (2 + \kappa_+) \right. \\
& + 2\nu (6153 + 5624 \kappa_+)] \} (nv)^5 - 12 \{ (60537 - 412820\nu + 216394\nu^2) \kappa_- + \delta [210 - 60537 \kappa_+ \\
& + 16294\nu^2 (2 + \kappa_+) + 2\nu (79897 + 34659 \kappa_+)] \} (nv)^3 v^2 + \{ 9 (11875 - 92497\nu + 57269\nu^2) \kappa_- \\
& + 3\delta [19862 - 35625 \kappa_+ + 44517\nu^2 (2 + \kappa_+) + 3\nu (88346 + 13763 \kappa_+)] \} (nv) v^4 \right) (v\Sigma) (nS) \\
& + \{ 30 [92054 - 51639\delta \kappa_- + 12\nu (-21737 + 2383\delta \kappa_- - 7117 \kappa_+) + 3159 \kappa_+ \\
& + 17664\nu^2 (2 + \kappa_+)] (nv)^6 - 30 [145074 - 60071\delta \kappa_- + 2\nu (-217150 + 17845\delta \kappa_- \\
& - 66746 \kappa_+) + 5839 \kappa_+ + 38456\nu^2 (2 + \kappa_+)] (nv)^4 v^2 + 6 [316786 - 73665\delta \kappa_- \\
& + 4\nu (-255289 + 13716\delta \kappa_- - 87213 \kappa_+) + 50505 \kappa_+ + 144912\nu^2 (2 + \kappa_+)] (nv)^2 v^4 \\
& - 6 [31638 - 1905\delta \kappa_- + \nu (-81524 + 2538\delta \kappa_- - 58068 \kappa_+) \\
& + 19521 \kappa_+ + 39720\nu^2 (2 + \kappa_+)] v^6 \} (nS)^2 \\
& + \left(90 \{ (9619 - 56962\nu + 27572\nu^2) \kappa_- + \delta [7744 - 9619 \kappa_+ + 616\nu^2 (2 + \kappa_+) \right. \\
& + 4\nu (-295 + 2812 \kappa_+)] \} (nv)^5 - 3 \{ 4 (60537 - 412820\nu + 216394\nu^2) \kappa_- \\
& + \delta [182127 - 242148 \kappa_+ + 65176\nu^2 (2 + \kappa_+) + \nu (58154 + 277272 \kappa_+)] \} (nv)^3 v^2 \\
& + \{ 9 (11875 - 92497\nu + 57269\nu^2) \kappa_- + 3\delta [35069 - 35625 \kappa_+ + 44517\nu^2 (2 + \kappa_+) \\
& + \nu (77244 + 41289 \kappa_+)] \} (nv) v^4 \right) (vS) (n\Sigma) + \{ -90 [-14140 - 9619\delta \kappa_- \\
& + \nu (45414 + 34105\delta \kappa_- - 53343 \kappa_+) - 2\nu^2 (7778 + 6739\delta \kappa_- - 17987 \kappa_+) + 9619 \kappa_+ \\
& + 1232\nu^3 (2 + \kappa_+)] (nv)^5 + 3 [-321287 - 242148\delta \kappa_- + \nu (1052063 + 964276\delta \kappa_- - 1448572 \kappa_+) \\
& + 242148 \kappa_+ + 130352\nu^3 (2 + \kappa_+) + \nu^2 (-17459 - 400200\delta \kappa_- + 954744 \kappa_+)] (nv)^3 v^2
\end{aligned}$$

$$\begin{aligned}
& -3 \left[-45701 - 35625 \delta \kappa_- + 15 \nu (9357 + 10626 \delta \kappa_- - 15376 \kappa_+) + 35625 \kappa_+ + 89034 \nu^3 (2 + \kappa_+) \right. \\
& \left. + \nu^2 (260391 - 63645 \delta \kappa_- + 146223 \kappa_+) \right] (nv) v^4 \} (v\Sigma) (n\Sigma) \\
& + \left(60 \{ -3 (9133 - 53426 \nu + 22008 \nu^2) \kappa_- \right. \\
& \left. + \delta [5136 + 27399 \kappa_+ + 8832 \nu^2 (2 + \kappa_+) - 4 \nu (29311 + 14250 \kappa_+)] \} (nv)^6 \right. \\
& - 60 \{ (-32955 + 204733 \nu - 90608 \nu^2) \kappa_- + \delta [19228 \nu^2 (2 + \kappa_+) + 5 (4094 + 6591 \kappa_+) \\
& - 3 \nu (70417 + 28197 \kappa_+)] \} (nv)^4 v^2 \\
& + 12 \{ -3 (20695 - 116396 \nu + 60728 \nu^2) \kappa_- + \delta [72796 + 62085 \kappa_+ \\
& + 72456 \nu^2 (2 + \kappa_+) - 6 \nu (88651 + 33643 \kappa_+)] \} (nv)^2 v^4 \\
& + \{ 36 (3571 - 11371 \nu + 8312 \nu^2) \kappa_- - 12 \delta [11822 + 10713 \kappa_+ + 19860 \nu^2 (2 + \kappa_+) \\
& - 13 \nu (3467 + 2331 \kappa_+)] \} v^6 \} (nS) (n\Sigma) \\
& + \{ -30 [20690 + 27399 \delta \kappa_- + 2 \nu^2 (-104335 + 18714 \delta \kappa_- - 75714 \kappa_+) \\
& - 27399 \kappa_+ + 17664 \nu^3 (2 + \kappa_+) + \nu (-46172 - 108639 \delta \kappa_- + 163437 \kappa_+)] \} (nv)^6 \\
& + 30 [12682 + 32955 \delta \kappa_- + 6 \nu^2 (-69356 + 9153 \delta \kappa_- - 37350 \kappa_+) - 32955 \kappa_+ \\
& + 38456 \nu^3 (2 + \kappa_+) + \nu (3736 - 144662 \delta \kappa_- + 210572 \kappa_+)] \} (nv)^4 v^2 \\
& - 6 [1658 + 62085 \delta \kappa_- + 6 \nu^2 (-187653 + 21220 \delta \kappa_- - 88506 \kappa_+) - 62085 \kappa_+ \\
& + 144912 \nu^3 (2 + \kappa_+) + \nu (159576 - 275523 \delta \kappa_- + 399693 \kappa_+)] \} (nv)^2 v^4 \\
& + 6 [2876 + 10713 \delta \kappa_- + 2 \nu^2 (-50069 + 11199 \delta \kappa_- - 41502 \kappa_+) - 10713 \kappa_+ \\
& + 39720 \nu^3 (2 + \kappa_+) + \nu (18530 - 32208 \delta \kappa_- + 53634 \kappa_+)] \} v^6 \} (n\Sigma)^2 \\
& + \{ -20 [18940 + 8466 \delta \kappa_- + 7008 \kappa_+ + 10680 \nu^2 (2 + \kappa_+) - 3 \nu (22146 + 2281 \delta \kappa_- \\
& + 10033 \kappa_+)] \} (nv)^6 + 12 [73878 + 24257 \delta \kappa_- + 23513 \kappa_+ + 44080 \nu^2 (2 + \kappa_+) \\
& - \nu (231290 + 20306 \delta \kappa_- + 101717 \kappa_+)] \} (nv)^4 v^2 - 12 [33624 \nu^2 (2 + \kappa_+) \\
& + 2 (29548 + 5583 \delta \kappa_- + 7800 \kappa_+) - \nu (160250 + 9813 \delta \kappa_- + 63483 \kappa_+)] \} (nv)^2 v^4 \\
& + 4 [48002 + 2721 \delta \kappa_- + 11397 \kappa_+ + 22296 \nu^2 (2 + \kappa_+) - \nu (113566 + 2568 \delta \kappa_- \\
& + 34857 \kappa_+)] \} v^6 \} \mathbf{S}^2 + \left(-20 \{ 6 (243 - 1768 \nu + 2782 \nu^2) \kappa_- + \delta [39345 - 1458 \kappa_+ \right. \\
& \left. + 10680 \nu^2 (2 + \kappa_+) - 38 \nu (3245 + 612 \kappa_+)] \} (nv)^6 + 12 \{ (744 - 15617 \nu + 37144 \nu^2) \kappa_- \right. \\
& \left. + \delta [134379 - 744 \kappa_+ + 44080 \nu^2 (2 + \kappa_+) - \nu (409370 + 81411 \kappa_+)] \} (nv)^4 v^2 \right. \\
& - 12 \{ 6 (-739 + 1501 \nu + 938 \nu^2) \kappa_- + \delta [88237 + 4434 \kappa_+ + 33624 \nu^2 (2 + \kappa_+) \\
& - 6 \nu (41937 + 8945 \kappa_+)] \} (nv)^2 v^4 + 4 \{ -3 (2892 - 7135 \nu + 4008 \nu^2) \kappa_- \\
& \left. + \delta [55499 + 8676 \kappa_+ + 22296 \nu^2 (2 + \kappa_+) - \nu (137578 + 32289 \kappa_+)] \} v^6 \} (S\Sigma) \right. \\
& \left. + \{ 20 [-23735 - 729 \delta \kappa_- + 729 \kappa_+ + 10680 \nu^3 (2 + \kappa_+) + \nu (85325 - 3162 \delta \kappa_- + 1704 \kappa_+) \right. \\
& - \nu^2 (181780 + 1503 \delta \kappa_- + 21753 \kappa_+)] \} (nv)^6 - 6 [2 \nu^2 (-586820 + 1734 \delta \kappa_- \\
& - 83145 \kappa_+) - 12 (11831 + 62 \delta \kappa_- - 62 \kappa_+) + 88160 \nu^3 (2 + \kappa_+) \\
& \left. + \nu (563466 - 32897 \delta \kappa_- + 31409 \kappa_+)] \} (nv)^4 v^2 + 12 [-35720 + 2217 \delta \kappa_- \right. \\
& \left. + \nu^2 (-340023 + 6999 \delta \kappa_- - 60669 \kappa_+) - 3 \nu (-56914 + 5223 \delta \kappa_- - 6701 \kappa_+) \right. \\
& \left. - 2217 \kappa_+ + 33624 \nu^3 (2 + \kappa_+)] \} (nv)^2 v^4 - 2 [2 \nu^2 (-160933 + 8580 \delta \kappa_- - 40869 \kappa_+) \right. \\
& \left. + 18 (-1493 + 482 \delta \kappa_- - 482 \kappa_+) + 44592 \nu^3 (2 + \kappa_+) \right. \\
& \left. + \nu (176176 - 26847 \delta \kappa_- + 44199 \kappa_+)] \} v^6 \} \mathbf{\Sigma}^2, \right. \\
f_8^1 = & \{ 4 [273931 - 90475 \delta \kappa_- + \nu (-101017 + 85054 \delta \kappa_- - 190731 \kappa_+) + 260591 \kappa_+ \\
& + 14 \nu^2 (2 + \kappa_+)] \} (nv)^2 + 4 [-89163 + 14451 \delta \kappa_- - 83310 \kappa_+ + 9122 \nu^2 (2 + \kappa_+) \\
& + \nu (-17229 - 13020 \delta \kappa_- + 14347 \kappa_+)] \} v^2 \} (vS)^2 \\
& + \left(4 \{ (-351066 + 637685 \nu - 340230 \nu^2) \kappa_- \right. \\
& \left. + \delta [-165796 + \nu (159708 - 275785 \kappa_+) + 351066 \kappa_+ + 14 \nu^2 (2 + \kappa_+)] \} (nv)^2 \right. \\
& + 4 \{ (97761 - 85171 \nu + 42958 \nu^2) \kappa_- + \delta [11030 - 97761 \kappa_+ + 9122 \nu^2 (2 + \kappa_+) \\
& + \nu (-92446 + 27367 \kappa_+)] \} v^2 \} (vS) (v\Sigma) + \{ -2 [570262 + 351066 \delta \kappa_- \\
& - 3 \nu (507492 + 152245 \delta \kappa_- - 386289 \kappa_+) + 2 \nu^2 (420201 + 85061 \delta \kappa_- - 360846 \kappa_+) \\
& - 351066 \kappa_+ + 28 \nu^3 (2 + \kappa_+)] \} (nv)^2 - 2 [-125332 - 97761 \delta \kappa_- \\
& + \nu (305810 + 56269 \delta \kappa_- - 251791 \kappa_+) - 2 \nu^2 (171541 + 8459 \delta \kappa_- - 35826 \kappa_+) \\
& + 97761 \kappa_+ + 18244 \nu^3 (2 + \kappa_+)] \} v^2 \} (v\Sigma)^2 + \{ 2 [-3056658 + 929364 \delta \kappa_- \\
& - 2218433 \kappa_+ + 46079 \nu^2 (2 + \kappa_+) + \nu (1676740 - 620106 \delta \kappa_- + 1491545 \kappa_+)] \} (nv)^3
\end{aligned}$$

$$\begin{aligned}
& + [3214028 - 874916 \delta \kappa_- + \nu (361732 + 302654 \delta \kappa_- - 80076 \kappa_+) + 2154322 \kappa_+ \\
& - 475778 \nu^2 (2 + \kappa_+)] (nv) v^2 \} (vS) (nS) + \left(\{ (3147797 - 5829107 \nu + 2434345 \nu^2) \kappa_- \right. \\
& + \delta [-801874 - 3147797 \kappa_+ + 46079 \nu^2 (2 + \kappa_+) + \nu (1009158 + 2111651 \kappa_+)] \} (nv)^3 \\
& + \{ (-1514619 + 1941197 \nu - 367419 \nu^2) \kappa_- + \delta [997862 + \nu (114558 - 191365 \kappa_+) \\
& + 1514619 \kappa_+ - 237889 \nu^2 (2 + \kappa_+)] \} (nv) v^2 \} (v\Sigma) (nS) \\
& + \{ 2 [5227548 - 1216896 \delta \kappa_- + \nu (-6893056 + 684456 \delta \kappa_- - 3387479 \kappa_+) \\
& + 3075029 \kappa_+ + 413377 \nu^2 (2 + \kappa_+)] (nv)^4 - 6 [1544838 - 328848 \delta \kappa_- \\
& + \nu (-2011336 + 138159 \delta \kappa_- - 853814 \kappa_+) + 907817 \kappa_+ + 82405 \nu^2 (2 + \kappa_+)] (nv)^2 v^2 \\
& + 4 [359181 - 57141 \delta \kappa_- + \nu (-616905 + 30384 \delta \kappa_- - 219799 \kappa_+) + 266628 \kappa_+ \\
& + 30484 \nu^2 (2 + \kappa_+)] v^4 \} (nS)^2 + \left(\{ (3147797 - 5829107 \nu + 2434345 \nu^2) \kappa_- \right. \\
& + \delta [-837370 - 3147797 \kappa_+ + 46079 \nu^2 (2 + \kappa_+) + \nu (125484 + 2111651 \kappa_+)] \} (nv)^3 \\
& + \{ (-1514619 + 1941197 \nu - 367419 \nu^2) \kappa_- + \delta [914702 + \nu (700104 - 191365 \kappa_+) \\
& + 1514619 \kappa_+ - 237889 \nu^2 (2 + \kappa_+)] \} (nv) v^2 \} (vS) (n\Sigma) + \{ [2673496 + 3147797 \delta \kappa_- \\
& + \nu^2 (1386252 + 1194133 \delta \kappa_- - 5417435 \kappa_+) - 3147797 \kappa_+ - 92158 \nu^3 (2 + \kappa_+) \\
& + \nu (-4318238 - 3970379 \delta \kappa_- + 10265973 \kappa_+)] (nv)^3 + [-688444 - 1514619 \delta \kappa_- \\
& + 3 \nu (-120938 + 355427 \delta \kappa_- - 1365173 \kappa_+) + 1514619 \kappa_+ + 475778 \nu^3 (2 + \kappa_+) \\
& + \nu^2 (-1396092 - 64765 \delta \kappa_- + 447495 \kappa_+)] (nv) v^2 \} (v\Sigma) (n\Sigma) \\
& + \left(2 \{ (-4291925 + 8939519 \nu - 3151201 \nu^2) \kappa_- + \delta [4026508 + 4291925 \kappa_+ \right. \\
& + 413377 \nu^2 (2 + \kappa_+) - 35 \nu (186627 + 116341 \kappa_+)] \} (nv)^4 \\
& - 6 \{ (-1236665 + 2307365 \nu - 635041 \nu^2) \kappa_- + \delta [1502404 + 1236665 \kappa_+ \\
& + 82405 \nu^2 (2 + \kappa_+) - \nu (2117145 + 991973 \kappa_+)] \} (nv)^2 v^2 \\
& + 4 \{ (-323769 + 478747 \nu - 152020 \nu^2) \kappa_- + \delta [429376 + 323769 \kappa_+ + 30484 \nu^2 (2 + \kappa_+) \\
& - \nu (690749 + 250183 \kappa_+)] \} v^4 \} (nS) (n\Sigma) + \{ [\nu (-6095254 + 6505727 \delta \kappa_- - 15089577 \kappa_+) \\
& - 5 (174488 + 858385 \delta \kappa_- - 858385 \kappa_+) - 826754 \nu^3 (2 + \kappa_+) \\
& + \nu^2 (12407868 - 1782289 \delta \kappa_- + 9926159 \kappa_+)] (nv)^4 + 3 [\nu^2 (-4525560 + 358723 \delta \kappa_- \\
& - 2342669 \kappa_+) + 85 (-980 + 14549 \delta \kappa_- - 14549 \kappa_+) + 164810 \nu^3 (2 + \kappa_+) \\
& + \nu (3418150 - 1649669 \delta \kappa_- + 4122999 \kappa_+)] (nv)^2 v^2 - 2 [-26330 + 323769 \delta \kappa_- \\
& + 2 \nu^2 (-777095 + 45626 \delta \kappa_- - 295809 \kappa_+) - 323769 \kappa_+ + 60968 \nu^3 (2 + \kappa_+) \\
& + \nu (1204762 - 364465 \delta \kappa_- + 1012003 \kappa_+)] v^4 \} (n\Sigma)^2 + \{ -4 [433968 - 47922 \delta \kappa_- \\
& + \nu (-715366 + 10725 \delta \kappa_- - 315989 \kappa_+) + 142766 \kappa_+ + 76576 \nu^2 (2 + \kappa_+)] (nv)^4 \\
& + 4 [645648 - 61356 \delta \kappa_- + \nu (-789234 + 15507 \delta \kappa_- - 356657 \kappa_+) + 187518 \kappa_+ \\
& + 80846 \nu^2 (2 + \kappa_+)] (nv)^2 v^2 - 8 [125822 - 7115 \delta \kappa_- + \nu (-74954 + 2894 \delta \kappa_- \\
& - 34242 \kappa_+) + 30553 \kappa_+ + 6601 \nu^2 (2 + \kappa_+)] v^4 \} \mathbf{S}^2 \\
& + \left(-8 \{ (-95344 + 259201 \nu - 59738 \nu^2) \kappa_- + \delta [354179 + 95344 \kappa_+ + 38288 \nu^2 (2 + \kappa_+) \right. \\
& - \nu (485847 + 163357 \kappa_+)] \} (nv)^4 + 8 \{ (-124437 + 308794 \nu - 71437 \nu^2) \kappa_- \\
& + \delta [510814 + 124437 \kappa_+ + 40423 \nu^2 (2 + \kappa_+) - \nu (564209 + 186082 \kappa_+)] \} (nv)^2 v^2 \\
& - 8 \{ (-37668 + 65596 \nu - 18177 \nu^2) \kappa_- + \delta [170713 + 37668 \kappa_+ + 6601 \nu^2 (2 + \kappa_+) \\
& - 22 \nu (5271 + 1688 \kappa_+)] \} v^4 \} (S\Sigma) + \left(4 [\nu^2 (-1257006 + 49013 \delta \kappa_- - 375727 \kappa_+) \right. \\
& + 16 (-12070 + 5959 \delta \kappa_- - 5959 \kappa_+) + 76576 \nu^3 (2 + \kappa_+) + \nu (1115375 - 211279 \delta \kappa_- \\
& + 401967 \kappa_+)] (nv)^4 - 4 [-298355 + 124437 \delta \kappa_- + \nu^2 (-1503589 + 55930 \delta \kappa_- - 428094 \kappa_+) \\
& - 124437 \kappa_+ + 80846 \nu^3 (2 + \kappa_+) + \nu (1600063 - 247438 \delta \kappa_- + 496312 \kappa_+)] (nv)^2 v^2 \\
& + 4 \{ -87541 + 37668 \delta \kappa_- - 6 \nu (-80620 + 8561 \delta \kappa_- - 21117 \kappa_+) - 37668 \kappa_+ \\
& + 13202 \nu^3 (2 + \kappa_+) + \nu^2 [12389 \delta \kappa_- - 9 (35739 + 9629 \kappa_+)] \} v^4 \} \mathbf{\Sigma}^2, \\
f_8^2 = & -2 [-902544 + 223199 \delta \kappa_- - 463094 \kappa_+ + 9198 \nu^2 (2 + \kappa_+) \\
& + \nu (-396690 - 754 \delta \kappa_- + 5787 \kappa_+)] (vS)^2 - 2 \{ (686293 - 899337 \nu - 6182 \nu^2) \kappa_- \\
& + \delta [-480890 - 686293 \kappa_+ + 9198 \nu^2 (2 + \kappa_+) + \nu (-419408 + 6541 \kappa_+)] \} (vS) (v\Sigma) \\
& + [-386486 - 686293 \delta \kappa_- + \nu (96445 + 452939 \delta \kappa_- - 1825525 \kappa_+) + 686293 \kappa_+ \\
& + 18396 \nu^3 (2 + \kappa_+) + \nu^2 (-910930 + 7690 \delta \kappa_- + 5392 \kappa_+)] (v\Sigma)^2 \\
& + [-6363300 + 2012530 \delta \kappa_- - 3332386 \kappa_+ + 140100 \nu^2 (2 + \kappa_+)
\end{aligned}$$

$$\begin{aligned}
& -2\nu(1950318 + 7900\delta\kappa_- + 169755\kappa_+)](nv)(vS)(nS) \\
& + \{ (2672458 - 3863205\nu - 38450\nu^2)\kappa_- + \delta[-2063276 - 2672458\kappa_+ \\
& + 70050\nu^2(2 + \kappa_+) - \nu(1928546 + 161855\kappa_+)] \} (nv)(vS)(nS) \\
& + \{ 2[7031031 - 1519937\delta\kappa_- + 2903969\kappa_+ - 3306\nu^2(2 + \kappa_+) \\
& + \nu(3359898 + 12907\delta\kappa_- + 727107\kappa_+)](nv)^2 + [-9180930 + 1353406\delta\kappa_- \\
& - 3337204\kappa_+ + 40356\nu^2(2 + \kappa_+) - 2\nu(1918032 + 4963\delta\kappa_- + 612954\kappa_+)]v^2 \} (nS)^2 \\
& + \{ (2672458 - 3863205\nu - 38450\nu^2)\kappa_- + \delta[70050\nu^2(2 + \kappa_+) \\
& - \nu(2054114 + 161855\kappa_+) - 2(908446 + 1336229\kappa_+)] \} (nv)(vS)(n\Sigma) \\
& + \{ 817592 + 2672458\delta\kappa_- - 2672458\kappa_+ - 140100\nu^3(2 + \kappa_+) \\
& + \nu(521708 - 1850675\delta\kappa_- + 7195591\kappa_+) + \nu^2[-54250\delta\kappa_- \\
& + 44(95009 + 8590\kappa_+)] \} (nv)(v\Sigma)(n\Sigma) + \{ -4(2211953 - 2682774\nu + 24161\nu^2)\kappa_- \\
& + 2\delta[6095633 + 4423906\kappa_+ - 3306\nu^2(2 + \kappa_+) + \nu(3369929 + 714200\kappa_+)] \} (nv)^2 \\
& + \{ -2(-2345305 + 2098821\nu + 326\nu^2)\kappa_- + 2\delta[-4610579 - 2345305\kappa_+ \\
& + 20178\nu^2(2 + \kappa_+) - \nu(1887473 + 607991\kappa_+)] \} v^2 \} (nS)(n\Sigma) \\
& + \{ [529810 - 4423906\delta\kappa_- + \nu(-9756929 + 2325674\delta\kappa_- - 11173486\kappa_+) \\
& + 4423906\kappa_+ + 6612\nu^3(2 + \kappa_+) - 2\nu^2(3364415 + 11254\delta\kappa_- + 702946\kappa_+)](nv)^2 \\
& + [5(-185396 + 469061\delta\kappa_- - 469061\kappa_+) - 40356\nu^3(2 + \kappa_+) \\
& + \nu^2(3647998 - 10252\delta\kappa_- + 1226234\kappa_+) + \nu(9257888 - 745415\delta\kappa_- \\
& + 5436025\kappa_+)]v^2 \} (n\Sigma)^2 + \{ -2[759660 - 171224\delta\kappa_- + 1669\nu\delta\kappa_- \\
& + 412592\kappa_+ + 22248\nu^2(2 + \kappa_+) + 24\nu(4621 + 8133\kappa_+)](nv)^2 \\
& + [1406688 - 302336\delta\kappa_- + 803672\kappa_+ - 7320\nu^2(2 + \kappa_+) + \nu(417852 + 2806\delta\kappa_- \\
& + 444750\kappa_+)]v^2 \} \mathbf{S}^2 + \{ -2\{(-583816 + 491373\nu - 28924\nu^2)\kappa_- \\
& + \delta[976976 + 583816\kappa_+ + 22248\nu^2(2 + \kappa_+) + \nu(80996 + 193523\kappa_+)] \} (nv)^2 \\
& - 4\{2(138251 - 95925\nu + 488\nu^2)\kappa_- + \delta[-499190 - 276502\kappa_+ + 1830\nu^2(2 + \kappa_+) \\
& - \nu(87269 + 110486\kappa_+)] \} v^2 \} (S\Sigma) + \{ [-118138 + 583816\delta\kappa_- \\
& - 583816\kappa_+ + 44496\nu^3(2 + \kappa_+) + \nu^2(34756 + 25586\delta\kappa_- + 361460\kappa_+) \\
& + \nu(2720885 - 148925\delta\kappa_- + 1316557\kappa_+)](nv)^2 + [357386 - 553004\delta\kappa_- \\
& + \nu(-2888077 + 81364\delta\kappa_- - 1187372\kappa_+) + \nu^2(-225662 + 854\delta\kappa_- - 442798\kappa_+) \\
& + 553004\kappa_+ + 7320\nu^3(2 + \kappa_+)]v^2 \} \mathbf{\Sigma}^2, \\
f_8^3 = & 3[11798 + 265\delta\kappa_- + 4\nu(-10242 + 463\delta\kappa_- - 5019\kappa_+) + 4977\kappa_+ + 672\nu^2(2 + \kappa_+)](nS)^2 \\
& + 6\{ -2(1178 - 5217\nu + 2020\nu^2)\kappa_- + \delta[4643 + 2356\kappa_+ + 336\nu^2(2 + \kappa_+) \\
& - 2\nu(10129 + 5482\kappa_+)] \} (nS)(n\Sigma) - 3[4\nu^2(-10066 + 547\delta\kappa_- - 6029\kappa_+) \\
& + 4(-2 + 589\delta\kappa_- - 589\kappa_+) + 672\nu^3(2 + \kappa_+) + \nu(10014 - 10699\delta\kappa_- + 15411\kappa_+)](n\Sigma)^2 \\
& + [-14234 - 265\delta\kappa_- - 4\nu(-9402 + 463\delta\kappa_- - 5019\kappa_+) - 4977\kappa_+ - 672\nu^2(2 + \kappa_+)] \mathbf{S}^2 \\
& + \{ 4(1178 - 5217\nu + 2020\nu^2)\kappa_- - 2\delta[3671 + 2356\kappa_+ + 336\nu^2(2 + \kappa_+) \\
& - 2\nu(9751 + 5482\kappa_+)] \} (S\Sigma) + [4\nu^2(-10234 + 547\delta\kappa_- - 6029\kappa_+) \\
& + 4(13 + 589\delta\kappa_- - 589\kappa_+) + 672\nu^3(2 + \kappa_+) + \nu(9942 - 10699\delta\kappa_- + 15411\kappa_+)] \mathbf{\Sigma}^2.
\end{aligned}$$

-
- [1] R. Abbott *et al.* (LIGO Scientific, VIRGO, KAGRA), GWTC-3: Compact Binary Coalescences Observed by LIGO and Virgo During the Second Part of the Third Observing Run, (2021), arXiv:2111.03606 [gr-qc].
- [2] A. H. Nitz, S. Kumar, Y.-F. Wang, S. Kastha, S. Wu, M. Schäfer, R. Dhurkunde, and C. D. Capano, 4-OGC: Catalog of gravitational waves from compact-binary mergers, (2021), arXiv:2112.06878 [astro-ph.HE].
- [3] S. Olsen, T. Venumadhav, J. Mushkin, J. Roulet, B. Zai-

-
- ckay, and M. Zaldarriaga, New binary black hole mergers in the LIGO–Virgo O3a data, (2022), arXiv:2201.02252 [astro-ph.HE].
- [4] P. Amaro-Seoane *et al.* (LISA), Laser Interferometer Space Antenna, (2017), arXiv:1702.00786 [astro-ph.IM].
- [5] M. Punturo *et al.*, The einstein telescope: A third-generation gravitational wave observatory, *Class. Quant. Grav.* **27**, 194002 (2010).
- [6] R. A. Porto, The Tune of Love and the Nature(ness) of Spacetime, *Fortsch. Phys.* **64**, 723 (2016), arXiv:1606.08895.
- [7] R. A. Porto, The Music of the Spheres: The Dawn of

- Gravitational Wave Science, (2017), arXiv:1703.06440 [physics.pop-ph].
- [8] M. Maggiore *et al.*, Science Case for the Einstein Telescope, JCAP **03**, 050, arXiv:1912.02622 [astro-ph.CO].
- [9] R. Alves Batista *et al.*, EuCAPT White Paper: Opportunities and Challenges for Theoretical Astroparticle Physics in the Next Decade, (2021), arXiv:2110.10074 [astro-ph.HE].
- [10] D. Kushnir, M. Zaldarriaga, J. A. Kollmeier, and R. Waldman, GW150914: Spin based constraints on the merger time of the progenitor system, Mon. Not. Roy. Astron. Soc. **462**, 844 (2016), arXiv:1605.03839 [astro-ph.HE].
- [11] W. M. Farr, S. Stevenson, M. Coleman Miller, I. Mandel, B. Farr, and A. Vecchio, Distinguishing Spin-Aligned and Isotropic Black Hole Populations With Gravitational Waves, Nature **548**, 426 (2017), arXiv:1706.01385 [astro-ph.HE].
- [12] J. Roulet, H. S. Chia, S. Olsen, L. Dai, T. Venumadhav, B. Zackay, and M. Zaldarriaga, Distribution of effective spins and masses of binary black holes from the LIGO and Virgo O1–O3a observing runs, Phys. Rev. D **104**, 083010 (2021), arXiv:2105.10580 [astro-ph.HE].
- [13] T. A. Callister, C.-J. Haster, K. K. Y. Ng, S. Vitale, and W. M. Farr, Who Ordered That? Unequal-mass Binary Black Hole Mergers Have Larger Effective Spins, Astrophys. J. Lett. **922**, L5 (2021), arXiv:2106.00521 [astro-ph.HE].
- [14] Y. Bouffanais, M. Mapelli, F. Santoliquido, N. Giacobbo, U. N. Di Carlo, S. Rastello, M. C. Artale, and G. Iorio, New insights on binary black hole formation channels after GWTC-2: young star clusters versus isolated binaries, Mon. Not. Roy. Astron. Soc. **507**, 5224 (2021), arXiv:2102.12495 [astro-ph.HE].
- [15] G. Franciolini, R. Cotesta, N. Loutrel, E. Berti, P. Pani, and A. Riotto, How to assess the primordial origin of single gravitational-wave events with mass, spin, eccentricity, and deformability measurements, (2021), arXiv:2112.10660 [astro-ph.CO].
- [16] A. Arvanitaki and S. Dubovsky, Exploring the String Axiverse with Precision Black Hole Physics, Phys. Rev. D **83**, 044026 (2011), arXiv:1004.3558 [hep-th].
- [17] D. Baumann, H. S. Chia, and R. A. Porto, Probing Ultralight Bosons with Binary Black Holes, Phys. Rev. D **99**, 044001 (2019), arXiv:1804.03208.
- [18] D. Baumann, H. S. Chia, R. A. Porto, and J. Stout, Gravitational Collider Physics, Phys. Rev. D **101**, 083019 (2020), arXiv:1912.04932.
- [19] S. Vitale, R. Lynch, J. Veitch, V. Raymond, and R. Sturani, Measuring the spin of black holes in binary systems using gravitational waves, Phys. Rev. Lett. **112**, 251101 (2014), arXiv:1403.0129 [gr-qc].
- [20] K. K. Ng, M. Isi, C.-J. Haster, and S. Vitale, Multiband gravitational-wave searches for ultralight bosons, Phys. Rev. D **102**, 083020 (2020), arXiv:2007.12793 [gr-qc].
- [21] R. Abbott *et al.* (LIGO Scientific, VIRGO, KAGRA), All-sky search for gravitational wave emission from scalar boson clouds around spinning black holes in LIGO O3 data, (2021), arXiv:2111.15507 [astro-ph.HE].
- [22] Q. Ding, X. Tong, and Y. Wang, Gravitational Collider Physics via Pulsar-Black Hole Binaries, Astrophys. J. **908**, 78 (2021), arXiv:2009.11106 [astro-ph.HE].
- [23] T. Takahashi and T. Tanaka, Axion clouds may survive the perturbative tidal interaction over the early inspiral phase of black hole binaries 10.1088/1475-7516/2021/10/031 (2021), arXiv:2106.08836 [gr-qc].
- [24] E. Maggio, P. Pani, and G. Raposo, Testing the nature of dark compact objects with gravitational waves, (2021), arXiv:2105.06410 [gr-qc].
- [25] P. Ajith *et al.*, The NINJA-2 catalog of hybrid post-Newtonian/numerical-relativity waveforms for non-precessing black-hole binaries, Class. Quant. Grav. **29**, 124001 (2012), [Erratum: Class.Quant.Grav. 30, 199401 (2013)], arXiv:1201.5319 [gr-qc].
- [26] B. Szilágyi, J. Blackman, A. Buonanno, A. Taracchini, H. P. Pfeiffer, M. A. Scheel, T. Chu, L. E. Kidder, and Y. Pan, Approaching the Post-Newtonian Regime with Numerical Relativity: A Compact-Object Binary Simulation Spanning 350 Gravitational-Wave Cycles, Phys. Rev. Lett. **115**, 031102 (2015), arXiv:1502.04953 [gr-qc].
- [27] T. Dietrich *et al.*, CoRe database of binary neutron star merger waveforms, Class. Quant. Grav. **35**, 24LT01 (2018), arXiv:1806.01625 [gr-qc].
- [28] L. Barack and A. Pound, Self-force and radiation reaction in general relativity, Rept. Prog. Phys. **82**, 016904 (2019), arXiv:1805.10385 [gr-qc].
- [29] L. Blanchet, Gravitational Radiation from Post-Newtonian Sources and Inspiral Compact Binaries, Living Rev. Rel. **17**, 2 (2014), arXiv:1310.1528 [gr-qc].
- [30] G. Schäfer and P. Jaranowski, Hamiltonian formulation of general relativity and post-Newtonian dynamics of compact binaries, Living Rev. Rel. **21**, 7 (2018), arXiv:1805.07240.
- [31] R. A. Porto, The effective field theorist’s approach to gravitational dynamics, Phys. Rept. **633**, 1 (2016), arXiv:1601.04914.
- [32] A. Buonanno, B. Iyer, E. Ochsner, Y. Pan, and B. S. Sathyaprakash, Comparison of post-Newtonian templates for compact binary inspiral signals in gravitational-wave detectors, Phys. Rev. D **80**, 084043 (2009), arXiv:0907.0700 [gr-qc].
- [33] M. Pürrer and C.-J. Haster, Gravitational waveform accuracy requirements for future ground-based detectors, Phys. Rev. Res. **2**, 023151 (2020), arXiv:1912.10055 [gr-qc].
- [34] D. Ferguson, K. Jani, P. Laguna, and D. Shoemaker, Assessing the readiness of numerical relativity for LISA and 3G detectors, Phys. Rev. D **104**, 044037 (2021), arXiv:2006.04272 [gr-qc].
- [35] S. Galadage, C. Talbot, T. Nagar, D. Jain, E. Thrane, and I. Mandel, Building Better Spin Models for Merging Binary Black Holes: Evidence for Nonspinning and Rapidly Spinning Nearly Aligned Subpopulations, Astrophys. J. Lett. **921**, L15 (2021), arXiv:2109.02424 [gr-qc].
- [36] T. Damour, P. Jaranowski, and G. Schäfer, Nonlocal-In-Time Action for the Fourth Post-Newtonian Conservative Dynamics of Two-Body Systems, Phys. Rev. D **89**, 064058 (2014), arXiv:1401.4548.
- [37] P. Jaranowski and G. Schäfer, Derivation of local-in-time fourth post-Newtonian ADM Hamiltonian for spinless compact binaries, Phys. Rev. **D92**, 124043 (2015), arXiv:1508.01016.
- [38] L. Bernard, L. Blanchet, A. Bohe, G. Faye, and S. Marsat, Fokker action of nonspinning compact binaries at the fourth post-Newtonian approximation, Phys. Rev. D **93**, 084037 (2016), arXiv:1512.02876.
- [39] L. Bernard, L. Blanchet, A. Bohe, G. Faye, and S. Marsat, Dimensional regularization of the IR diver-

- gences in the Fokker action of point-particle binaries at the fourth post-Newtonian order, *Phys. Rev.* **D96**, 104043 (2017), arXiv:1706.08480.
- [40] T. Marchand, L. Bernard, L. Blanchet, and G. Faye, Ambiguity-Free Completion of the Equations of Motion of Compact Binary Systems at the Fourth Post-Newtonian Order, *Phys. Rev. D* **97**, 044023 (2018), arXiv:1707.09289.
- [41] D. Bini, T. Damour, and A. Geralico, Novel approach to binary dynamics: application to the fifth post-Newtonian level, *Phys. Rev. Lett.* **123**, 231104 (2019), arXiv:1909.02375 [gr-qc].
- [42] C. Galley, A. Leibovich, R. A. Porto, and A. Ross, Tail Effect in Gravitational Radiation Reaction: Time Nonlocality and Renormalization Group Evolution, *Phys. Rev. D* **93**, 124010 (2016), arXiv:1511.07379.
- [43] R. A. Porto and I. Rothstein, Apparent Ambiguities in the Post-Newtonian Expansion for Binary Systems, *Phys. Rev. D* **96**, 024062 (2017), arXiv:1703.06433.
- [44] R. A. Porto, Lamb Shift and the Gravitational Binding Energy for Binary Black Holes, *Phys. Rev. D* **96**, 024063 (2017), arXiv:1703.06434.
- [45] S. Foffa and R. Sturani, Dynamics of the gravitational two-body problem at fourth post-Newtonian order and at quadratic order in the Newton constant, *Phys. Rev. D* **87**, 064011 (2013), arXiv:1206.7087 [gr-qc].
- [46] S. Foffa, R. A. Porto, I. Rothstein, and R. Sturani, Conservative dynamics of binary systems to fourth Post-Newtonian order in the EFT approach II: Renormalized Lagrangian, *Phys. Rev.* **D100**, 024048 (2019), arXiv:1903.05118.
- [47] T. Marchand, Q. Henry, F. Larrouturou, S. Marsat, G. Faye, and L. Blanchet, The mass quadrupole moment of compact binary systems at the fourth post-Newtonian order, *Class. Quant. Grav.* **37**, 215006 (2020), arXiv:2003.13672 [gr-qc].
- [48] Q. Henry, G. Faye, and L. Blanchet, The current-type quadrupole moment and gravitational-wave mode $(\ell, m) = (2, 1)$ of compact binary systems at the third post-Newtonian order, *Class. Quant. Grav.* **38**, 185004 (2021), arXiv:2105.10876 [gr-qc].
- [49] J. Blümlein, A. Maier, P. Marquard, and G. Schäfer, The fifth-order post-Newtonian Hamiltonian dynamics of two-body systems from an effective field theory approach: potential contributions, *Nucl. Phys. B* **965**, 115352 (2021), arXiv:2010.13672 [gr-qc].
- [50] J. Blümlein, A. Maier, P. Marquard, and G. Schäfer, The fifth-order post-Newtonian Hamiltonian dynamics of two-body systems from an effective field theory approach, (2021), arXiv:2110.13822 [gr-qc].
- [51] S. Foffa and R. Sturani, Hereditary terms at next-to-leading order in two-body gravitational dynamics, *Phys. Rev. D* **101**, 064033 (2020), arXiv:1907.02869 [gr-qc].
- [52] G. L. Almeida, S. Foffa, and R. Sturani, Tail contributions to gravitational conservative dynamics, (2021), arXiv:2110.14146 [gr-qc].
- [53] J. Blümlein, A. Maier, P. Marquard, and G. Schäfer, Testing binary dynamics in gravity at the sixth post-Newtonian level, *Phys. Lett. B* **807**, 135496 (2020), arXiv:2003.07145 [gr-qc].
- [54] J. Blümlein, A. Maier, P. Marquard, and G. Schäfer, The 6th post-Newtonian potential terms at $O(G_N^4)$, *Phys. Lett. B* **816**, 136260 (2021), arXiv:2101.08630 [gr-qc].
- [55] D. Bini, T. Damour, and A. Geralico, Sixth post-Newtonian local-in-time dynamics of binary systems, *Phys. Rev. D* **102**, 024061 (2020), arXiv:2004.05407 [gr-qc].
- [56] D. Bini, T. Damour, and A. Geralico, Sixth post-Newtonian nonlocal-in-time dynamics of binary systems, *Phys. Rev. D* **102**, 084047 (2020), arXiv:2007.11239 [gr-qc].
- [57] S. Marsat, A. Bohe, G. Faye, and L. Blanchet, Next-to-next-to-leading order spin-orbit effects in the equations of motion of compact binary systems, *Class. Quant. Grav.* **30**, 055007 (2013), arXiv:1210.4143 [gr-qc].
- [58] A. Bohé, S. Marsat, and L. Blanchet, Next-to-next-to-leading order spin-orbit effects in the gravitational wave flux and orbital phasing of compact binaries, *Class. Quant. Grav.* **30**, 135009 (2013), arXiv:1303.7412 [gr-qc].
- [59] R. A. Porto, Post-Newtonian Corrections to the Motion of Spinning Bodies in NRGR, *Phys. Rev. D* **73**, 104031 (2006), arXiv:gr-qc/0511061.
- [60] R. A. Porto and I. Z. Rothstein, The hyperfine Einstein-Infeld-Hoffmann potential, *Phys. Rev. Lett.* **97**, 021101 (2006), arXiv:gr-qc/0604099 [gr-qc].
- [61] R. A. Porto and I. Rothstein, Spin(1)Spin(2) Effects in the Motion of Inspiralling Compact Binaries at Third Order in the Post-Newtonian Expansion, *Phys. Rev. D* **78**, 044012 (2008), [Erratum: *Phys. Rev. D* **81** (2010) 029904], arXiv:0802.0720.
- [62] R. A. Porto and I. Z. Rothstein, Next to Leading Order Spin(1)Spin(1) Effects in the Motion of Inspiralling Compact Binaries, *Phys. Rev.* **D78**, 044013 (2008), arXiv:0804.0260.
- [63] R. A. Porto, A. Ross, and I. Z. Rothstein, Spin induced multipole moments for the gravitational wave flux from binary inspirals to third Post-Newtonian order, *JCAP* **1103**, 009, arXiv:1007.1312 [gr-qc].
- [64] R. A. Porto, A. Ross, and I. Z. Rothstein, Spin induced multipole moments for the gravitational wave amplitude from binary inspirals to 2.5 Post-Newtonian order, *JCAP* **1209**, 028, arXiv:1203.2962 [gr-qc].
- [65] A. Bohé, Alejandro, G. Faye, S. Marsat, and E. K. Porter, Quadratic-in-spin effects in the orbital dynamics and gravitational-wave energy flux of compact binaries at the 3PN order, *Class. Quant. Grav.* **32**, 195010 (2015), arXiv:1501.01529 [gr-qc].
- [66] G. Cho, B. Pardo, and R. A. Porto, Gravitational radiation from inspiralling compact objects: Spin-spin effects completed at the next-to-leading post-Newtonian order, *Phys. Rev. D* **104**, 024037 (2021), arXiv:2103.14612 [gr-qc].
- [67] W. D. Goldberger and I. Z. Rothstein, An Effective field theory of gravity for extended objects, *Phys. Rev.* **D73**, 104029 (2006), arXiv:hep-th/0409156.
- [68] W. D. Goldberger and A. Ross, Gravitational radiative corrections from effective field theory, *Phys. Rev.* **D81**, 124015 (2010), arXiv:0912.4254.
- [69] R. A. Porto, Next-to-Leading Order Spin-Orbit Effects in the Motion of Inspiralling Compact Binaries, *Class. Quant. Grav.* **27**, 205001 (2010), arXiv:1005.5730.
- [70] I. Z. Rothstein, TASI lectures on effective field theories, (2003), arXiv:hep-ph/0308266 [hep-ph].
- [71] S. Marsat, A. Bohé, L. Blanchet, and A. Buonanno, Next-to-leading tail-induced spin-orbit effects in the gravitational radiation flux of compact binaries, *Class. Quant. Grav.* **31**, 025023 (2014), arXiv:1307.6793 [gr-qc].
- [72] M. Levi and J. Steinhoff, Complete conservative dy-

- namics for inspiralling compact binaries with spins at the fourth post-Newtonian order, JCAP **09**, 029, arXiv:1607.04252 [gr-qc].
- [73] A. Antonelli, C. Kavanagh, M. Khalil, J. Steinhoff, and J. Vines, Gravitational spin-orbit coupling through third-subleading post-Newtonian order: from first-order self-force to arbitrary mass ratios, Phys. Rev. Lett. **125**, 011103 (2020), arXiv:2003.11391 [gr-qc].
- [74] A. Antonelli, C. Kavanagh, M. Khalil, J. Steinhoff, and J. Vines, Gravitational spin-orbit and aligned spin₁-spin₂ couplings through third-subleading post-Newtonian orders, Phys. Rev. D **102**, 124024 (2020), arXiv:2010.02018 [gr-qc].
- [75] M. Khalil, Gravitational spin-orbit dynamics at the fifth-and-a-half post-Newtonian order, Phys. Rev. D **104**, 124015 (2021), arXiv:2110.12813 [gr-qc].
- [76] Z. Liu, R. A. Porto, and Z. Yang, Spin Effects in the Effective Field Theory Approach to Post-Minkowskian Conservative Dynamics, JHEP **06**, 012, arXiv:2102.10059 [hep-th].
- [77] D. Kosmopoulos and A. Luna, Quadratic-in-spin Hamiltonian at $\mathcal{O}(G^2)$ from scattering amplitudes, JHEP **07**, 037, arXiv:2102.10137 [hep-th].
- [78] H. Tagoshi, M. Shibata, T. Tanaka, and M. Sasaki, Post-Newtonian expansion of gravitational waves from a particle in circular orbits around a rotating black hole: Up to $\mathcal{O}(v^8)$ beyond the quadrupole formula, Phys. Rev. D **54**, 1439 (1996), arXiv:gr-qc/9603028.
- [79] E. Poisson, Absorption of mass and angular momentum by a black hole: Time-domain formalisms for gravitational perturbations, and the small-hole / slow-motion approximation, Phys. Rev. D **70**, 084044 (2004), arXiv:gr-qc/0407050.
- [80] R. A. Porto, Absorption Effects due to Spin in the World-line Approach to Black Hole Dynamics, Phys. Rev. D **77**, 064026 (2008), arXiv:0710.5150.
- [81] W. D. Goldberger, J. Li, and I. Z. Rothstein, Non-conservative effects on spinning black holes from world-line effective field theory, JHEP **06**, 053, arXiv:2012.14869 [hep-th].
- [82] N. T. Maia, C. R. Galley, A. K. Leibovich, and R. A. Porto, Radiation reaction for spinning bodies in effective field theory I: Spin-orbit effects, Phys. Rev. **D96**, 084064 (2017), arXiv:1705.07934.
- [83] N. T. Maia, C. R. Galley, A. K. Leibovich, and R. A. Porto, Radiation reaction for spinning bodies in effective field theory II: Spin-spin effects, Phys. Rev. **D96**, 084065 (2017), arXiv:1705.07938.
- [84] S. Marsat, Cubic order spin effects in the dynamics and gravitational wave energy flux of compact object binaries, Class. Quant. Grav. **32**, 085008 (2015), arXiv:1411.4118 [gr-qc].
- [85] J. M. Martín-García *et al.*, xact: Efficient tensor computer algebra for the wolfram language, www.xact.es (2019).