

DEUTSCHE ELEKTRONEN-SYCHROTRON DESY

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1920-1921. The first year of the new century was a period of great change and development for the country. The economy was growing rapidly, and there was a sense of optimism and progress.

The year began with a general election, which resulted in a hung parliament. This led to a coalition government between the Conservative Party and the Liberal Party, under the leadership of Prime Minister David Lloyd George.

One of the key issues of the year was the question of Home Rule for Ireland. The Irish Parliamentary Party had been pushing for Home Rule since the late 1800s, and finally achieved it in 1914.

Another important event was the start of World War I. The war began in August 1914, and Britain entered the conflict on the side of France and Russia.

The year also saw the introduction of the Representation of the People Act, which gave women the right to vote for the first time.

Overall, 1920-1921 was a year of significant political and social change in the United Kingdom.

1922-1923. The year began with a general election, which resulted in a Conservative government under Prime Minister Andrew Bonar Law.

One of the key issues of the year was the question of Ulster Unionism, which sought to maintain Northern Ireland's status as part of the United Kingdom.

The year also saw the introduction of the National Insurance Act, which provided a safety net for the poor and unemployed.

Overall, 1922-1923 was a year of political stability and social reform in the United Kingdom.

1924-1925. The year began with a general election, which resulted in a Labour government under Prime Minister Ramsay MacDonald.

One of the key issues of the year was the question of the British Empire, particularly the relationship between Britain and India.

The year also saw the introduction of the Trade Descriptions Act, which protected consumers from misleading advertising.

Overall, 1924-1925 was a year of political change and social reform in the United Kingdom.

1926-1927. The year began with a general election, which resulted in a Conservative government under Prime Minister Stanley Baldwin.

One of the key issues of the year was the question of the British Empire, particularly the relationship between Britain and India.

The year also saw the introduction of the Health and Safety at Work Act, which protected workers from workplace accidents.

Overall, 1926-1927 was a year of political stability and social reform in the United Kingdom.

1928-1929. The year began with a general election, which resulted in a Conservative government under Prime Minister Neville Chamberlain.

One of the key issues of the year was the question of the British Empire, particularly the relationship between Britain and India.

The year also saw the introduction of the National Health Service, which provided free medical care for all citizens.

Overall, 1928-1929 was a year of political stability and social reform in the United Kingdom.

*Observation of F Meson Production in High Energy  $e^+e^-$  Annihilation*

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Submitted to Physics letters

November 1983

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##### Supported by the Konrad Adenauer Stiftung, on leave from the University of  
Science and Technology of China, Hefei  
\$ Supported by the Deutsches Bundesministerium für Forschung und Techno-  
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## Supported by the UK Science and Engineering Research Council

### Supported by the Minerva Gesellschaft für Forschung mbH

#### Supported by the US Department of Energy contract DE-AC02-76ER00881

Nonstrange charmed mesons have been extensively studied(1-5), while little is known about those carrying strangeness(6-8). Recently, the CLEO group (9) has reported the observation of the F meson produced by  $e^+e^-$  annihilation at a c.m. energy of  $W=10.5$  GeV. The F meson has been detected in the  $\varphi\pi$  decay mode with a mass of  $1.970 \pm 0.005 \pm 0.005$  GeV. While this value is not in disagreement with that obtained in the first measurement(6) of the F,  $m_F = 2.03 \pm 0.06$  GeV, it is at variance with the photoproduction results of the F which, averaged over all their measured decay modes, yield  $m_F = 2.02 \pm 0.01$  GeV (7-8). We present in this letter data on F production in  $e^+e^-$  annihilation which confirm the CLEO result at high energies.

The experiment was performed with the TASSO detector (10) at the storage ring PETRA at DESY. Data were taken at c.m. energies  $W=14, 22$  and  $25$  GeV with an integrated luminosity of  $5.4 pb^{-1}$  and between  $30$  and  $42$  GeV with an average  $\bar{W}=34.7$  GeV and an integrated luminosity of  $77.6 pb^{-1}$ . A total number of  $29186$  hadronic events were selected using the same cuts as in ref.11. A momentum resolution for charged particles of  $\sigma_p/p = 0.01/(2.9+p^2)^{1/2}$  was obtained by using the average beam position as a constraint in the track reconstruction(4).

F production was searched for in the decay mode  $F^\pm \rightarrow \varphi\pi^\pm \rightarrow K^+K^-\pi^\pm$ . Particle identification was used to reduce the number of wrong mass assignments. For  $|\cos\vartheta| \leq 0.8$ , where  $\vartheta$  is the angle with respect to the beam direction, particle identification was done by time-of-flight(TOF) counters (12) placed between the central drift chamber and the solenoid at a radius of  $1.32$  m. The  $\pi/K$  separation with the TOF counters is possible for momenta  $p \leq 1.0$  GeV/c. Since most of the  $\varphi$ 's coming from high momentum F mesons are expected to be energetic, the kaons were not required to be positively identified by the TOF counters.

Candidates for  $\varphi$  mesons were searched for by looking at invariant mass combinations of oppositely charged tracks interpreted as kaons. However, particles positively identified as pions or protons by the TOF counters were not considered as kaons<sup>1</sup>. The  $K^+K^-$  mass spectrum is shown in Fig.1 for all mass combinations and for the combinations with a fractional energy  $x_{KK} = 2E_{KK}/W \geq 0.4$  (shaded histogram). The latter shows a clear  $\varphi$  signal. A background (parametrized in terms of the  $K^+$  momentum in the  $K^+K^-$  rest frame multiplying a polynomial) plus a gaussian shape for the  $\varphi$  was fitted to the shaded histogram, yielding a mass were not required to be positively identified by the TOF counters.

of  $1023 \pm 2$  MeV in agreement with the standard  $\varphi$  mass,  $m_\varphi = 1.020$  GeV (13), and a width (FWHM) of  $21 \pm 7$  MeV, which is consistent with the width expected from our detector resolution. We define as  $\varphi$ -candidates all  $K^+K^-$  mass combinations within  $\pm 15$  MeV of the nominal  $\varphi$  mass.

The  $\varphi$  candidates were combined with any of the other charged tracks in the event interpreted as a pion ( $\pi$  particles identified as sure kaons or protons in the TOF system were not taken). The resulting  $K^+K^-\pi^\pm$  mass spectrum is shown in Fig. 2a for combinations with a fractional energy  $x_{KK} = 2E_{KK}/W > 0.4$ . A clear enhancement is seen at a mass around  $1.97$  GeV. No signal is observed when  $K^+K^-$  combinations in the  $\varphi$  control region ( $1.05 \leq M_{KK} \leq 1.10$  GeV) are selected (Fig. 2b). From this and the fact that the  $\varphi$  control region is close to the  $\varphi$  mass we conclude that the  $1.97$  GeV signal observed in Fig. 2a is associated with the  $\varphi\pi$  channel. The same conclusion can be drawn from a comparison of the  $K^+K^-$  mass distribution for the  $K^+K^-\pi^\pm$  mass combination in the  $1.97$  GeV region and in a control region.

In order to determine the mass and width of the observed signal, a gaussian plus a background term of the form  $dN/dM = a_0 + a_1 M + a_2 M^2$  were fitted to the  $KK\pi$  mass spectrum shown in Fig.2a. The fit yielded a signal of  $49 \pm 14$  events and a mass of  $1.975 \pm 0.009$  GeV. The systematic uncertainty of the mass calibration in this range was found to be less than  $10$  MeV. This was determined from  $K^0_s \rightarrow \pi^+\pi^-$  decays, from the position of the  $\varphi$  peak, and from a measurement of the total c.m. energy with  $\mu^+\mu^-$  pairs. The width of the gaussian ( $0.064 \pm 0.021$  GeV FWHM) is consistent with the expected mass resolution and hence with a small natural width. The properties of the  $1.97$  GeV signal (mass, width and  $\varphi\pi$  decay mode) agree with those of the  $F(1.97)$  observed by the CLEO group(9).

We now consider F production at  $W > 30$  GeV ( $\bar{W}=34.7$  GeV). For the determination of the cross sections, the number of F mesons was found from the  $K^+K^-\pi^\pm$  mass distributions for  $x_{KK}$  intervals between  $0.3$  and  $1.0$  using fits with a polynomial background plus a gaussian for the F. Monte Carlo events of the type  $e^+e^- \rightarrow c\bar{c}$  including gluon emission and radiative effects and leading to  $F^\pm \rightarrow \varphi\pi^\pm$  were generated<sup>2</sup> and passed through the detector and the event selection criteria(11). These events were used to compute the detector acceptance and efficiency. The overall detection efficiency for the chain  $F^\pm \rightarrow \varphi\pi^\pm \rightarrow K^+K^-$  was found to be  $0.48$  at  $x_{KK} = 0.30$  and to decrease to  $0.26$  at  $x_{KK} = 0.9$ , the average being  $0.31$ . The data were also corrected for the other decay modes of the  $\varphi$ , assuming a branching ratio  $B(\varphi \rightarrow K^+K^-) = 0.49 \pm 0.01$  (13).

In Fig.3 we present the  $F^\pm$  scaled cross-section multiplied by the branching ratio,  $s/\beta \cdot d\sigma/dx \cdot B(F^\pm \rightarrow \varphi\pi^\pm)$ , ( $s=W^2$ ,  $\beta=p_F/E_F$ ) as a function of the scaled F energy  $x = 2E_F/W$ . The errors shown include only the statistical uncertainties. There is an  $\sum W_i = 1.0$ . Using this, we called a particle

<sup>1</sup> For each particle type i ( $i=\pi, k, p$ ) the probabilities  $W_i$  were calculated according to  $W_i = C f_i(p) \exp(-(r_m - r_i)^2 / 2 \sigma_r^2)$  where  $r_m$  is the measured TOF,  $r_i$  the expected TOF for hypothesis i,  $\sigma_r$  being the resolution for the time measurement,  $f_i(p)$  the momentum dependent particle fraction for the hypothesis according to ref.12 and C a normalization constant such that  $\sum W_i = 1.0$ .

a sure pion, if  $W_\pi \geq 0.9$  for  $0.3 \leq p_\pi \leq 1.0$  GeV/c

a sure kaon, if  $W_K \geq 0.6$  and  $0.1 \leq M \leq 0.8$  GeV for  $0.3 \leq p_K \leq 1.0$  GeV/c

a sure proton, if  $W_p \geq 0.6$  and  $0.6 \leq p_p \leq 1.4$  GeV/c,  
where M is the mass of the particle as calculated from the momentum and TOF measurements.

<sup>2</sup> From the spin properties of the F and the  $\varphi$ , an isotropic distribution was assumed for the  $\pi^\pm$  in the  $\varphi\pi$  rest frame. The decay  $\varphi \rightarrow K^+K^-$  was generated with a cosine $\vartheta_K$  distribution, where  $\vartheta_K$  is the angle between the  $K^+$  in the  $\varphi$  rest frame and the  $\varphi$  direction in the  $\varphi\pi$  rest frame.

overall systematic error of  $\pm 30\%$  due to uncertainties in the detection efficiency and in the background determination. The shape of the distribution is similar to our  $D^*$  cross section (4) and can be described well by the parametrization of ref.14:

$$\frac{s}{\beta} \cdot \frac{d\sigma}{dx} \sim \frac{1}{x \left[ 1 - \frac{1}{x} - \frac{\varepsilon}{1-x} \right]^2} \quad (1)$$

A fit to the data of Fig.3 yielded  $\varepsilon=0.45\pm 0.25$ , which can be compared with our value of  $\varepsilon=0.18\pm 0.07$  obtained for  $D^*$  production(4).

The integrated measured  $F^\pm$  cross section for  $x \geq 0.3$  yielded, relative to the  $\mu$  pair cross section ( $\sigma_{\mu\mu} = 4\pi\alpha^2/3s = 0.072\text{nb}$  at  $34.7\text{GeV}$ ),  $R_F(x \geq 0.3)B(F^\pm \rightarrow \varphi\pi^\pm) = \sigma_F(x > 0.3)B(F^\pm \rightarrow \varphi\pi^\pm)/\sigma_{\mu\mu} = 0.061\pm 0.012\pm 0.018$ . Using the parametrization given in eq.(1) to extrapolate to  $x \leq 0.3$ , the total  $F^\pm \rightarrow \varphi\pi^\pm$  production was found to be  $R_F B(F^\pm \rightarrow \varphi\pi^\pm) = 0.064\pm 0.013\pm 0.019$ . This can be compared with the expected total yield for primary strange charmed meson production,  $R(\bar{c}\bar{s}) = 2(\star/\star)4.010.167 = 0.486$ , where  $\star/\star$  is the relative contribution of the c quark to the total cross-section, 4.01 is our measured R value (11) and 0.167 is the assumed probability to pick up an  $s\bar{s}$  pair from the sea(15). From the observed  $F^\pm \rightarrow \varphi\pi^\pm$  yield and under the above assumptions we then find for the branching ratio  $B(F^\pm \rightarrow \varphi\pi^\pm) = 0.13\pm 0.03\pm 0.04$ , where the systematic error includes only the uncertainties in the measured cross section. This value is larger than the value of  $B(F^\pm \rightarrow \varphi\pi^\pm) \approx 0.044$  obtained in ref.9 in a similar way. Note, however, that our result for  $B(F^\pm \rightarrow \varphi\pi^\pm)$  may be reduced by as much as 25% due to possible F's originating from primary bottom quarks at low x.

In summary we have observed F meson production in the  $\varphi\pi^\pm$  final state with a mass of  $1.975\pm 0.009\pm 0.010$  GeV and a width consistent with our mass resolution. The F production yield times the branching ratio was found to be  $R_F(x \geq 0.3)B(F^\pm \rightarrow \varphi\pi^\pm) = 0.061\pm 0.012\pm 0.018$ .

We gratefully acknowledge the support of the DESY directorate and the PETRA machine group. We wish to thank the technical service groups at DESY and all the engineers and technicians at the collaborating institutes for their valuable help. One of us (Y.E.) wishes to thank the Minna-James-Heinemann Stiftung of Hannover for partial support. Those of us from outside DESY wish to thank the DESY directorate for the hospitality extended to them while working at DESY.

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## Figure captions

- Fig 1: Invariant mass distribution of oppositely charged particles (interpreted as kaons). Data from all c.m. energies are included. The shaded histogram (note different scale at the right hand side) is for all combinations with  $x_{KK} \geq 0.4$ . The curve shows the fit of a gaussian plus a background term (see text).
- Fig 2: Distributions of the  $K^+K^-$  mass with  $x_{KK} \geq 0.4$ . Data from all c.m. energies are included. The curves show the result of the fits as described in the text.
  - a) For  $K^+K^-$  combinations in the  $\varphi$  region ( $1.005 < M(K^+K^-) < 1.035$  GeV). The curve shows the fit of a gaussian plus a background term (see text).
  - b) For  $K^+K^-$  combinations in the  $\varphi$  control region ( $1.05 < M(K^+K^-) < 1.10$  GeV). The curve is a fit to a smooth background.
- Fig 3 : The scaled cross sections  $s/\beta \cdot d\sigma/dx \cdot B(F^\pm \rightarrow \varphi\pi^\pm)$  for F meson production at  $\bar{W}=34.7$  GeV. The curve shown is the result of the fit according to eq.(1).

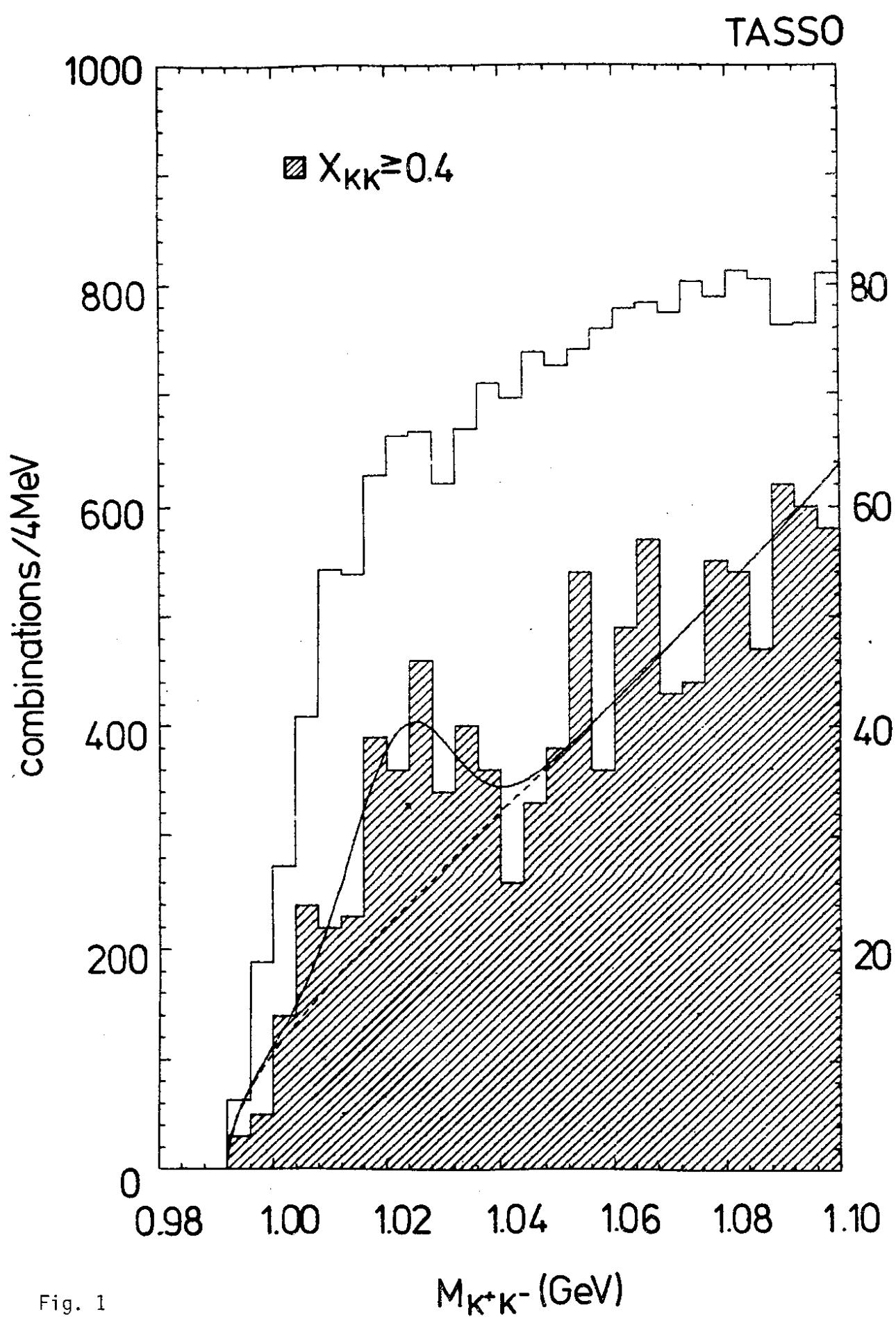


Fig. 1

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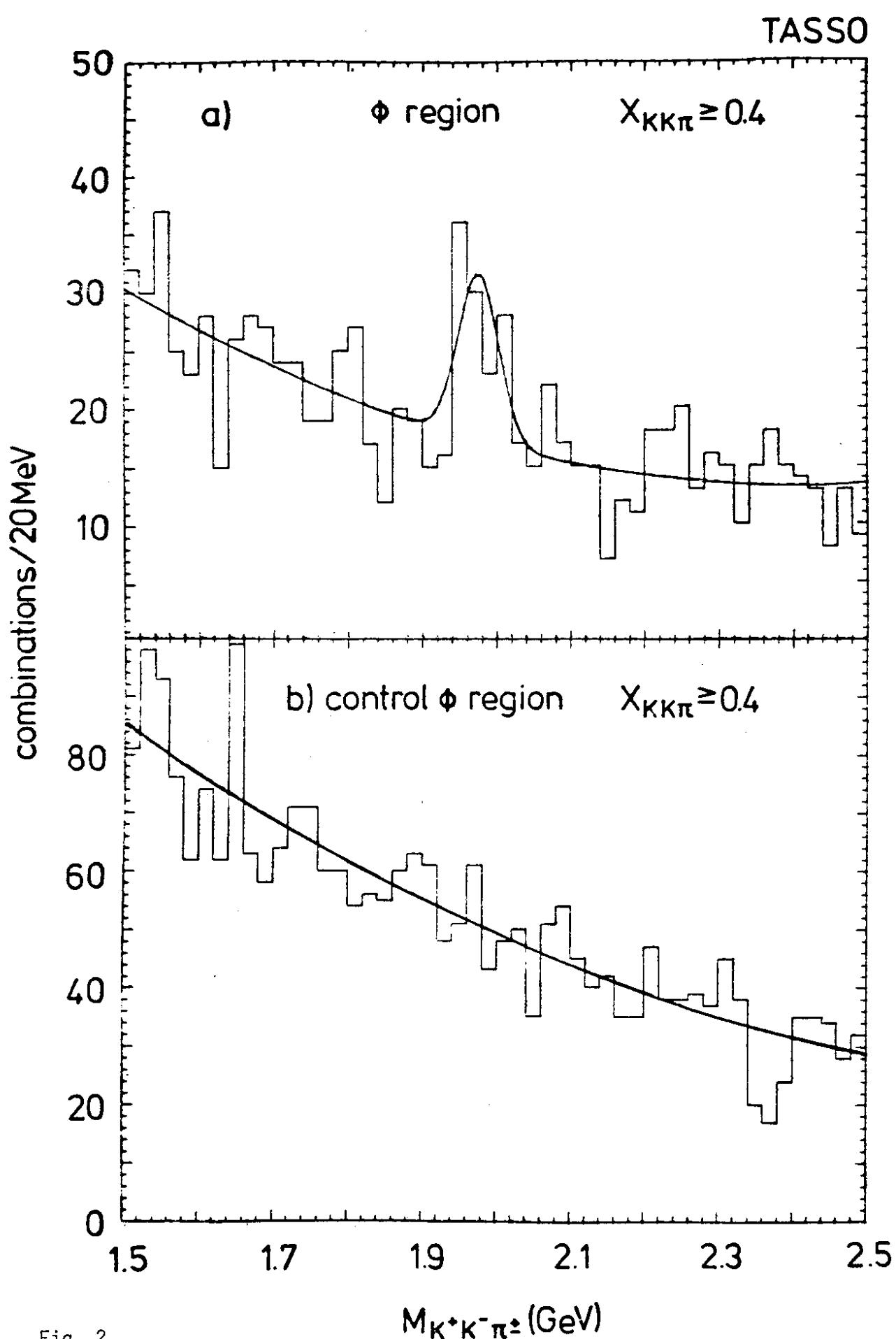


Fig. 2

$M_{K^+ K^- \pi^\pm}$  (GeV)

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TASSO

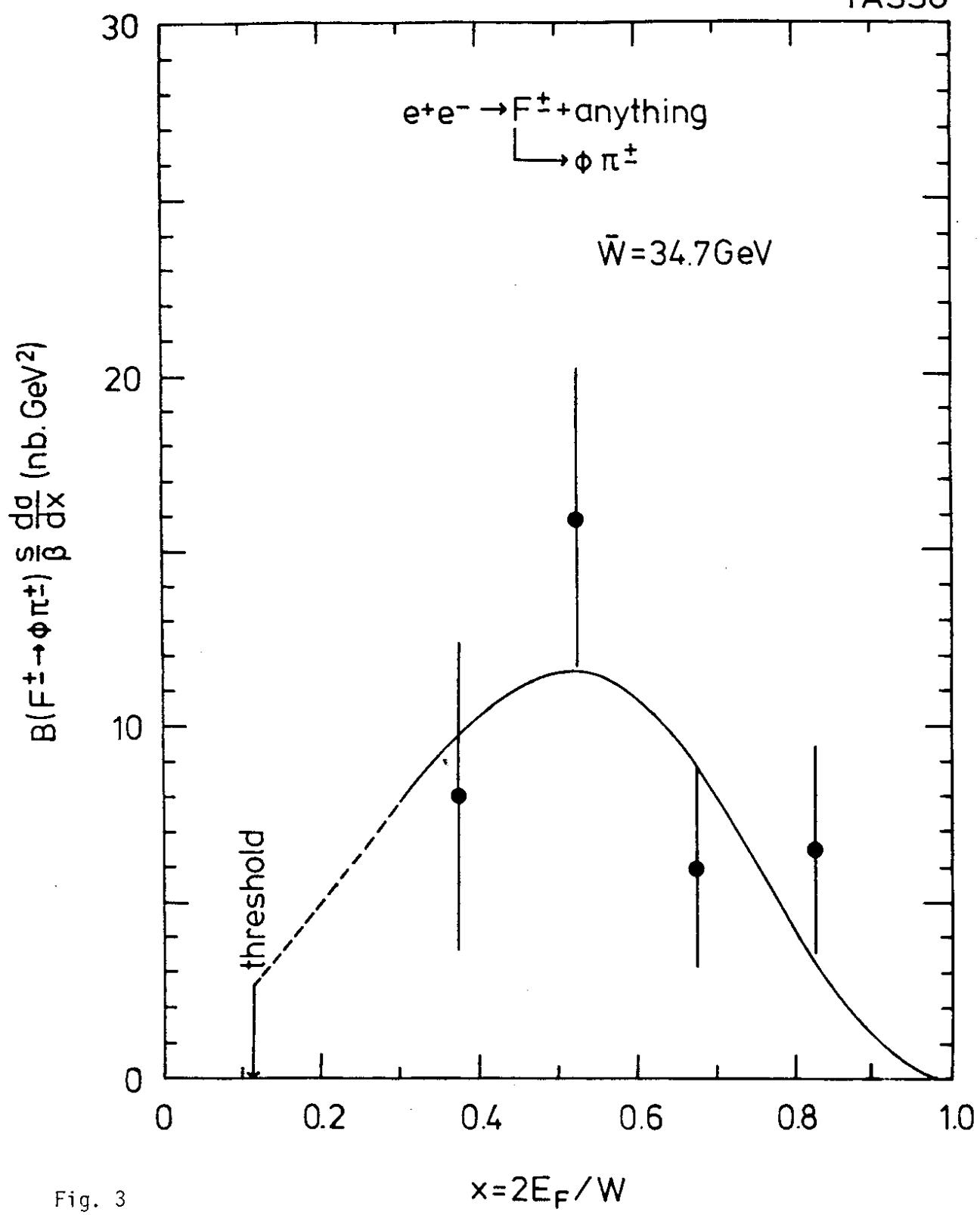


Fig. 3